



**Sharif University of Technology  
School of Mechanical Engineering  
Center of Excellence in Energy Conversion**

## **Advanced Thermodynamics**

### **Lecture 19**

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$\emptyset$  Monoatomic ideal gas:

- $\emptyset$  They only have  $e_e$  and  $e_t$ .
- $\emptyset$   $e_e$  is at the base level for them, they are non-degenerate.
- $\emptyset$  The electronic ground level degeneracy,  $g_{e0}$ , may be subtracted from the total partition functions.

$$Z = g_{e0} \sum_j g_{jt} e^{-\frac{e_{jt}}{KT}} = g_{e0} Z_t$$

Translational partition function

$$\bar{S} = \bar{S}_t + \bar{S}_{e0} = \bar{R} \left[ \ln \left( \frac{Z_t}{N} \right) + 1 \right] + \frac{U_t}{T} + \bar{R} \ln g_{e0}$$

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$$Z_t = \sum_{\text{levels}} g_j e^{-\frac{e_j}{KT}} \quad \text{or} \quad Z = \sum_{\text{quantum states}} e^{-\frac{e_j}{KT}}$$

Ø The energy of each quantum state due to Schrodinger Wave Eq.:

$$e_j = \frac{h^2}{8m} \left( \frac{K_x^2}{a^2} + \frac{K_y^2}{b^2} + \frac{K_z^2}{c^2} \right)$$

Ø  $K_x$ ,  $K_y$ , and  $K_z$  are quantum numbers

$$Z_t = \sum_{K_x=1}^{\infty} \sum_{K_y=1}^{\infty} \sum_{K_z=1}^{\infty} e^{-\left(\frac{h^2}{8mKT}\right) \left( \frac{K_x^2}{a^2} + \frac{K_y^2}{b^2} + \frac{K_z^2}{c^2} \right)}$$

$K_x$ ,  $K_y$ , and  $K_z$  are such huge numbers  $\rightarrow$

$$Z_t = \left[ \int_0^{\infty} e^{-\left(\frac{h^2}{8mKT}\right) \left( \frac{K_x^2}{a^2} \right)} dK_x \right] \left[ \int_0^{\infty} e^{-\left(\frac{h^2}{8mKT}\right) \left( \frac{K_y^2}{b^2} \right)} dK_y \right] \left[ \int_0^{\infty} e^{-\left(\frac{h^2}{8mKT}\right) \left( \frac{K_z^2}{c^2} \right)} dK_z \right]$$

$$\int_0^{\infty} e^{-x^2} dx = \frac{1}{2} \sqrt{\frac{p}{a}} \Rightarrow Z_t = V \left( \frac{2pmKT}{h^2} \right)^{\frac{3}{2}}$$

$$\ln Z_t = \ln V + \frac{3}{2} \ln T + \frac{3}{2} \ln \left( \frac{2pmK}{h^2} \right)$$

Ø Pressure of monoatomic ideal gas

$$P = NKT \left( \frac{\partial \ln Z}{\partial V} \right)_T = \frac{NKT}{V}$$

Ø Energy of monoatomic ideal gas

$$U = NKT^2 \left( \frac{\partial \ln Z}{\partial T} \right)_V = \frac{3}{2} NKT$$

Ø Entropy of monoatomic ideal gas

$$S = NKT^2 \left( \frac{\partial \ln Z}{\partial T} \right)_V = NK \left[ \frac{3}{2} \ln T + \ln \left( \frac{V}{N} \right) + \frac{3}{2} \ln \left( \frac{2pmK}{h^2} \right) + \frac{5}{2} \right]$$

$$\text{For one mole} \rightarrow \bar{s} = \frac{3}{2} \bar{R} \ln T + \bar{R} \ln v + \frac{\bar{R}}{N_A} \left[ \frac{3}{2} \ln \left( \frac{2pmK}{h^2} \right) + \frac{5}{2} \right]$$

$$\Rightarrow \boxed{\bar{s} = \bar{c}_v \ln T + \bar{R} \ln v + s_0}$$

For monoatomic ideal gas:

$$\ln Z_t = \ln V + \ln C + \frac{3}{2} \ln T$$

$$\left( \frac{\partial \ln Z_t}{\partial T} \right)_V = \frac{3}{2} T$$

$$\bar{U}_t = \frac{3}{2} \bar{R}T, \bar{h}_t = \frac{5}{2} \bar{R}T$$

$$\bar{c}_{V0} = \bar{c}_{Vt} = \frac{3}{2} \bar{R}, \bar{c}_{P0} = \bar{c}_{Pt} = \frac{5}{2} \bar{R}$$

$$\bar{s}_t = \bar{R} \left[ \ln \left( \frac{Z_t}{N} \right) + 1 \right] + \frac{\bar{U}_t}{T} = \bar{R} \left[ \ln \left( \frac{Z_t}{N} \right) + \frac{5}{2} \right]$$

- Ø In addition to  $s$ ,  $s_0$  will be calculated in statistical thermodynamics.

Ø Find Entropy of monoatomic Nitrogen at  $25^{\circ}C$  and  $0.1 \text{ MPa}$

Given:  $M = 14.0067$  and  $g_{e_0} = 4$

For monoatomic ideal gas:

$$\frac{Z_t}{N} = 2.5944 \frac{M^{\frac{3}{2}} T [K]^{\frac{5}{2}}}{P [KPa]} = 2.087 \times 10^6$$

$$\bar{s}_t = \bar{R} \left[ \ln \left( \frac{Z_t}{N} \right) + 1 \right] + \frac{\bar{U}_t}{T} = \bar{R} \left[ \ln \left( \frac{Z_t}{N} \right) + \frac{5}{2} \right] = 141.73 \frac{\text{KJ}}{\text{Kmole K}}$$

$$\bar{s}_{e_0} = \bar{R} \ln g_{e_0} = 11.526 \frac{\text{KJ}}{\text{Kmole K}}$$

$$\rightarrow \boxed{s = \bar{s}_t + \bar{s}_{e_0} = 153.298 \frac{\text{KJ}}{\text{Kmole K}}}$$

$\emptyset$  The validity of  $e^a \gg 1 \rightarrow g_j \gg N_j$

For He:

$$e^a = \frac{Z}{N} = \frac{Z_t + 1}{N} = 2.5944 \frac{M^{3/2} T [K]^{5/2}}{P [KPa]} = 2.5944 \frac{(4.0076)^{3/2} T^{5/2}}{100}$$

He ia an inert gas ( $g_{e_0} = 1$ ).

T=5K (around C.P.)

$$\rightarrow e^a = 1.1014 \text{ (It's better not to use C.B.M.)}$$

T=50K (cryogenic temp.)

$$\rightarrow e^a = 367.26 \text{ (C.B.M ??)}$$

T=500K

$$\rightarrow e^a = 116155.64 \text{ (C.B.M. is properly valid at this condition)}$$

$\emptyset$  Find  $N_{j \text{ M.P.}}$  for the Fermi-Dirac distribution function.

$$W_{F.D.} = \prod_j \left[ \frac{g_j !}{(g_j - N_j) ! N_j !} \right] \rightarrow \ln W = \sum_j [\ln g_j ! - \ln(g_j - N_j) ! - \ln N_j !]$$

$$\Rightarrow \ln W = \sum_j [g_j \ln g_j - (g_j - N_j) \ln(g_j - N_j) - N_j \ln N_j]$$

At fixed U, V, and N  $\rightarrow$

$$d \ln W = \sum_j [dN_j + \ln(g_j - N_j) dN_j - dN_j - \ln N_j dN_j] =$$

$$\sum_j \left[ -\ln \left( \frac{g_j - N_j}{N_j} \right) + a + b e_j \right] dN_j = 0 \rightarrow$$

$$N_{j \text{ M.P.}} = \frac{g_j}{e^a e^{b e_j} + 1}$$

( for B.M. =  $\frac{g_j}{e^a e^{b e_j}}$  and for B.E. =  $\frac{g_j}{e^a e^{b e_j} - 1}$  )