



Sharif Quantum Information
Group

Topological Quantum Computation



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New Advances in Condensed Matter Physics:
Quantum transport, topological effects and energy conversion in low-dimensional systems

Khiva, Uzbekistan, 2017





Objectives

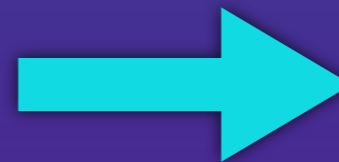
- To understand the basic ideas of:

Topological Qubit

Topological Order

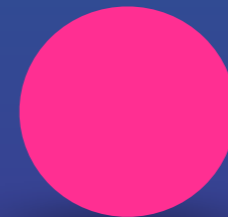
Kitaev Model

Topological Quantum Computation





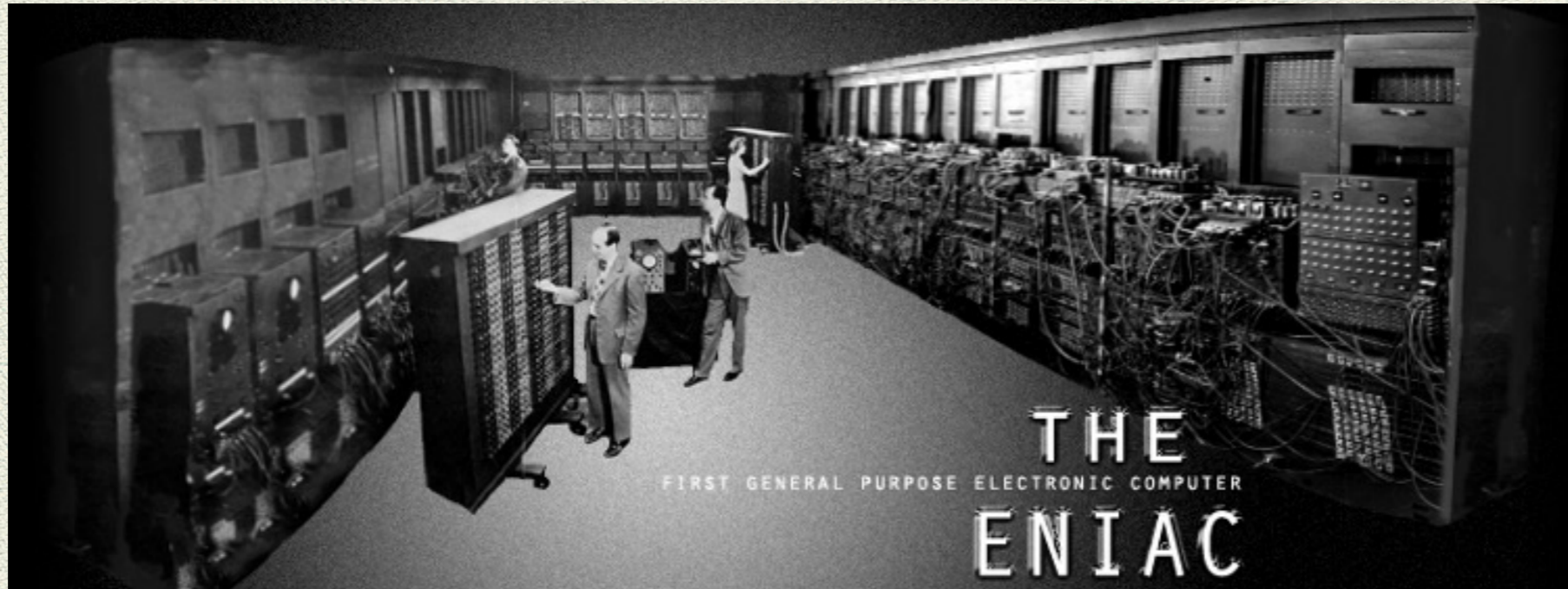
Classical Bits



Quantum Bits

$$|\text{blue/red}\rangle = a |\text{blue}\rangle + b |\text{red}\rangle$$

The early computers and computes of today



Intel Broadwell-EP Xeon

7.2 Billion Transistors



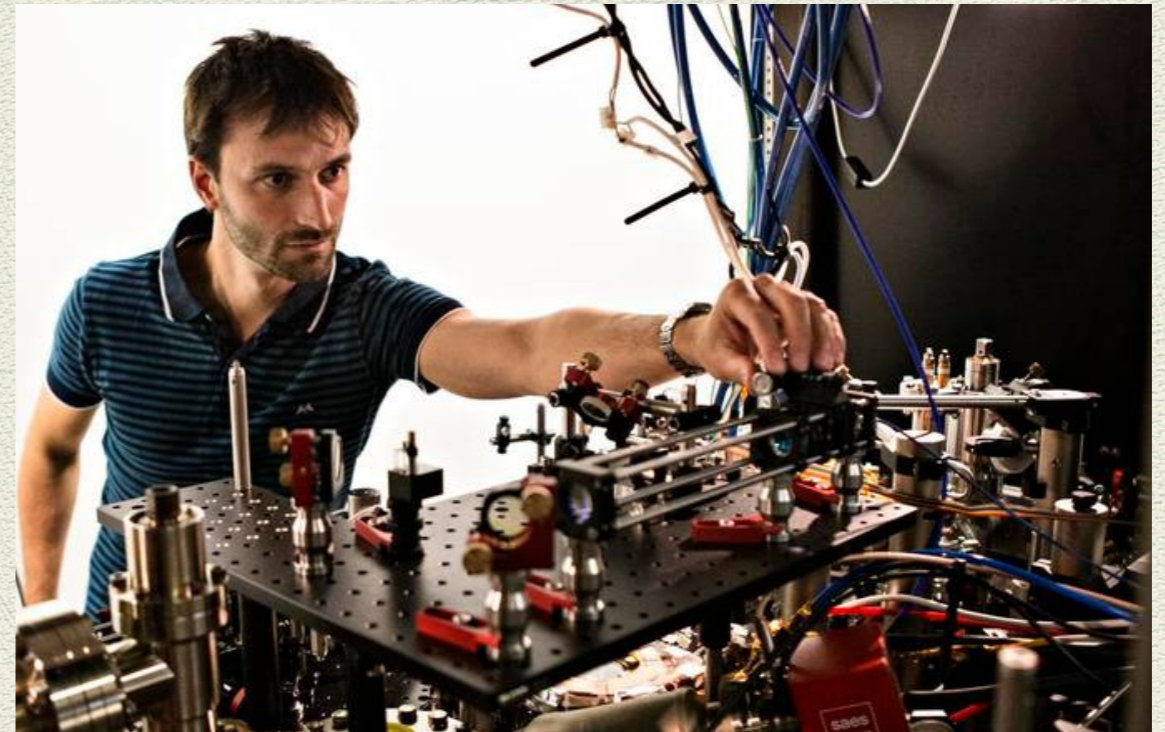
Computers of today and those of tomorrow



16 Giga Bytes

3.8 GHz

1300 US dollars



1 Kilo byte

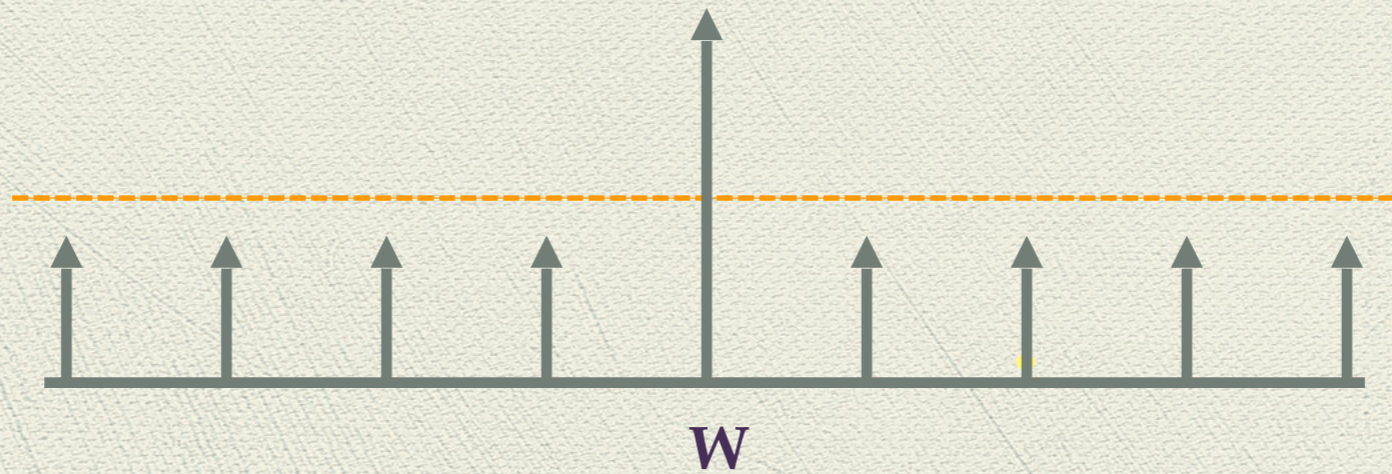
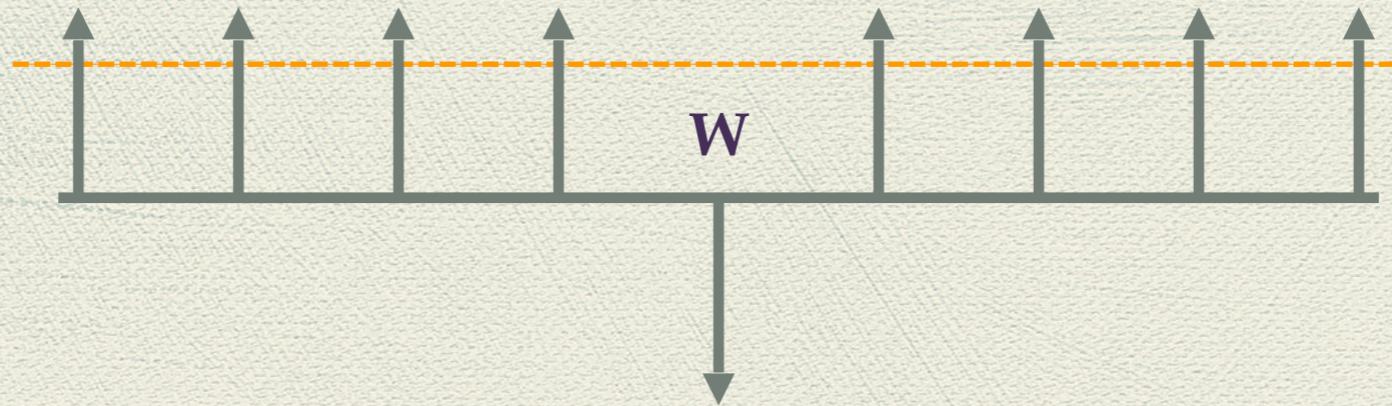
1 KHz

10s of millions of dollars

$$\sum_{x=1}^N |x\rangle$$



$$\sum_{x=1}^N (-1)^{f(x)} |x\rangle$$



Classical Bits Have

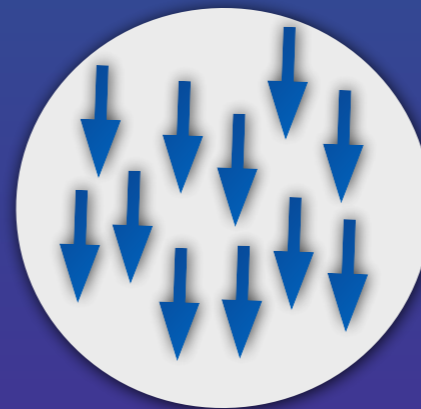
Four

Very Good Properties

1- Bits are Macroscopic Objects

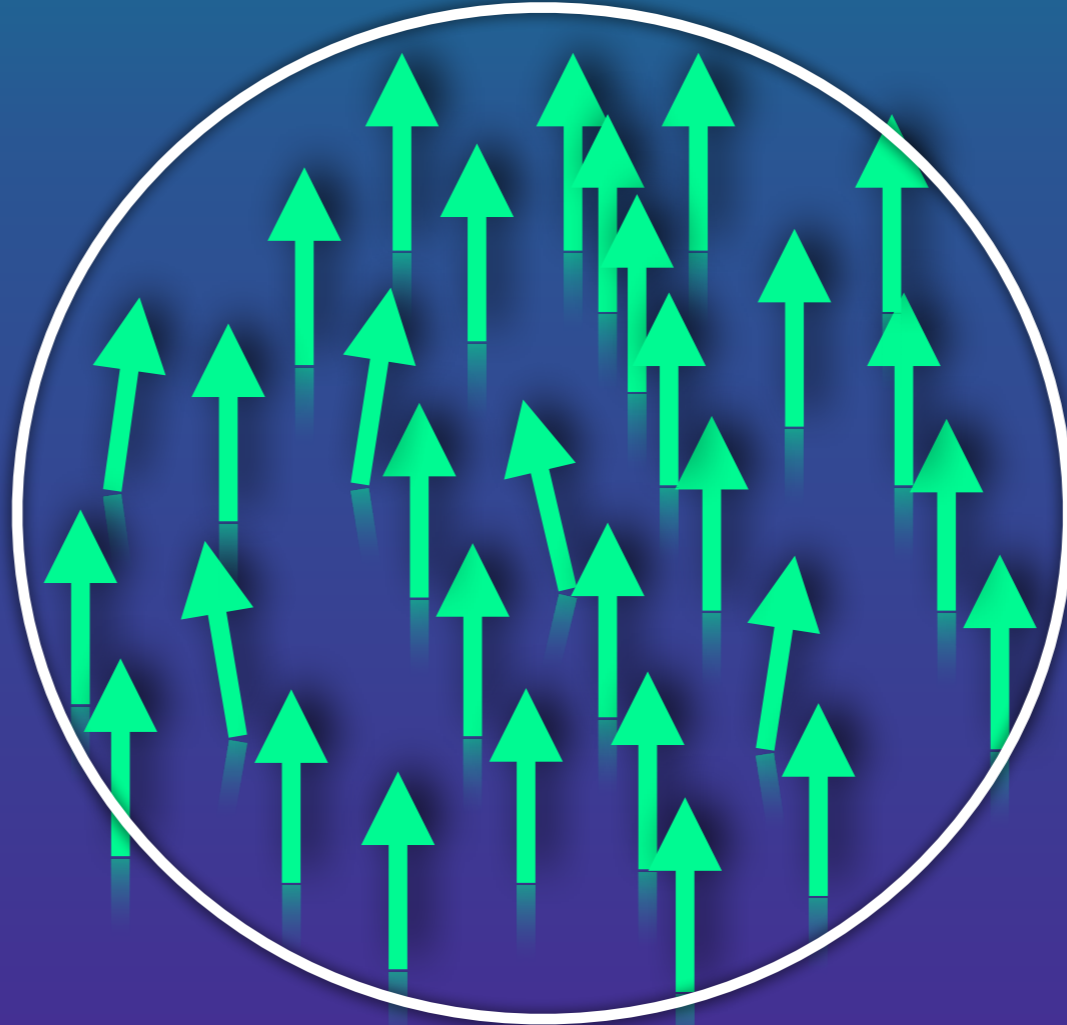


0

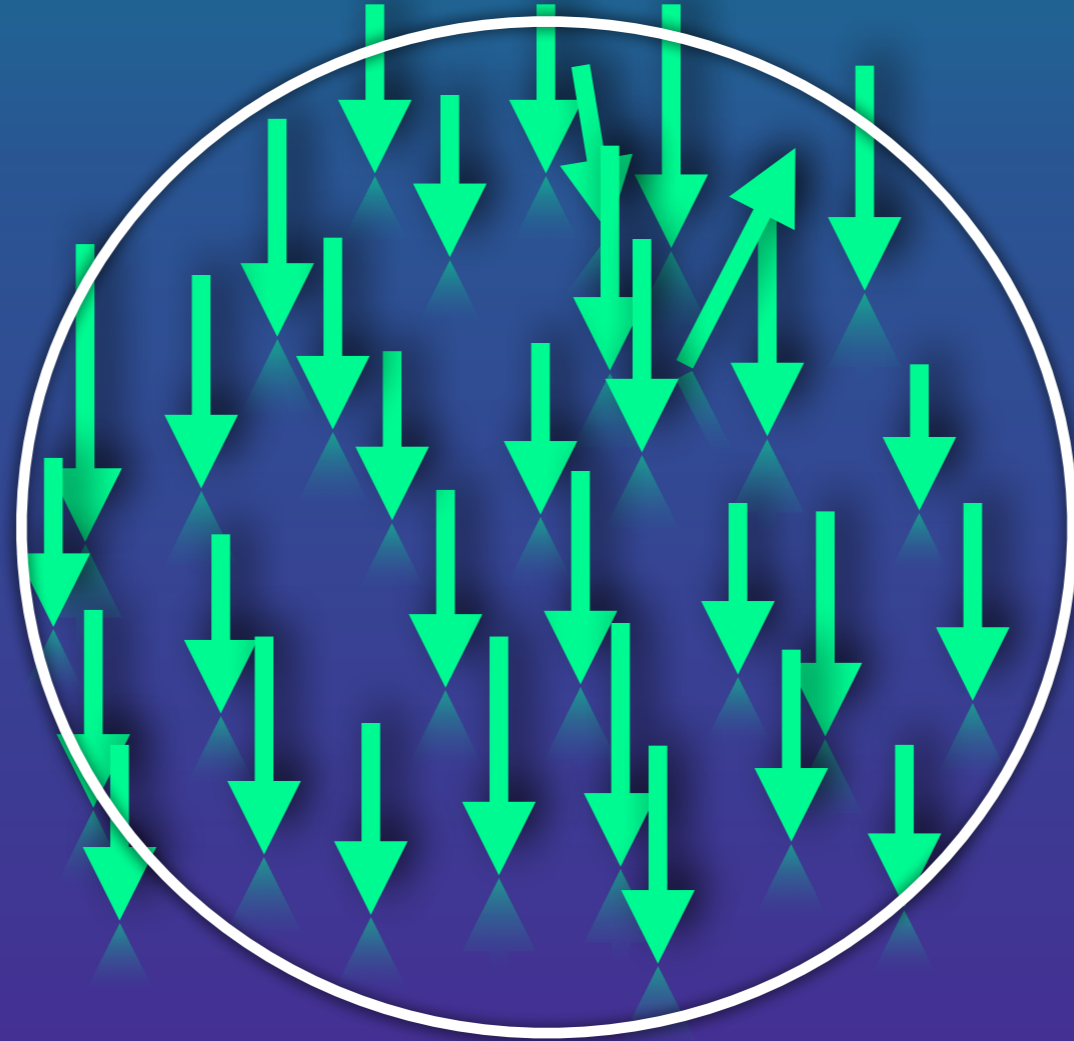


1

0

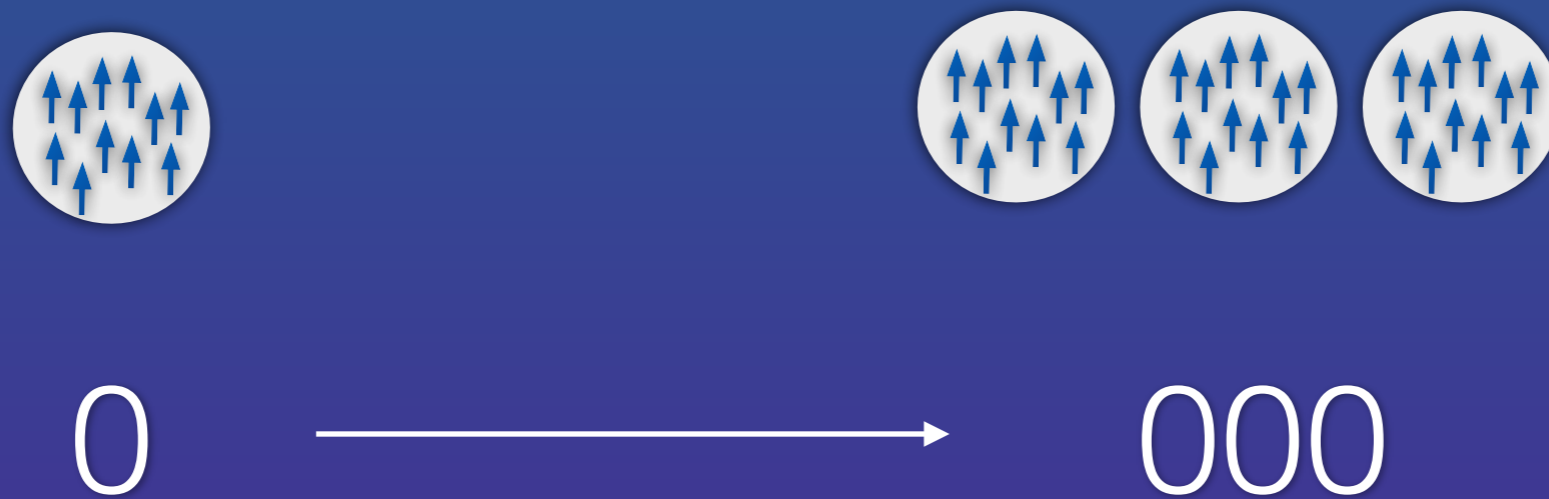


1



13

2- Bits can be cloned



3- Errors are discrete



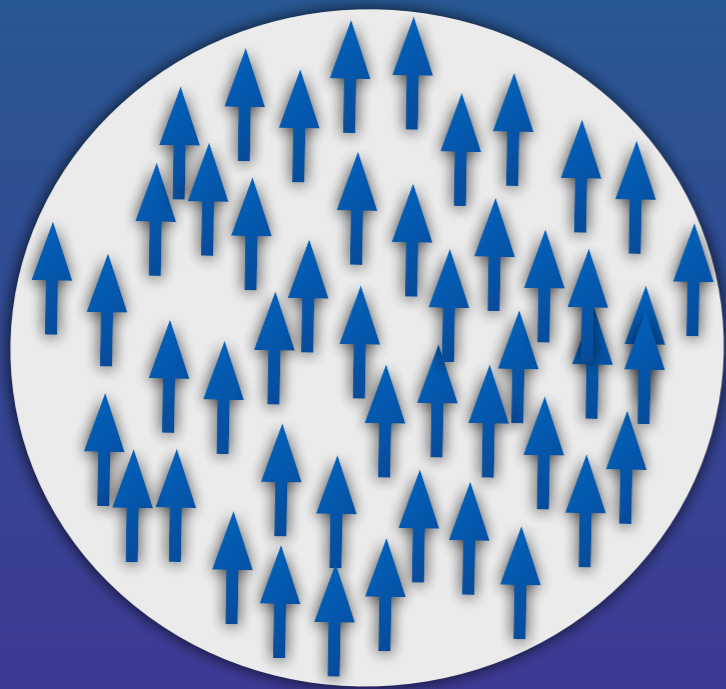
4- Bits can be observed

010 → 000

And corrected

Qubits have exactly the
opposite properties

They are microscopic



Bit



Qubit

They cannot be cloned

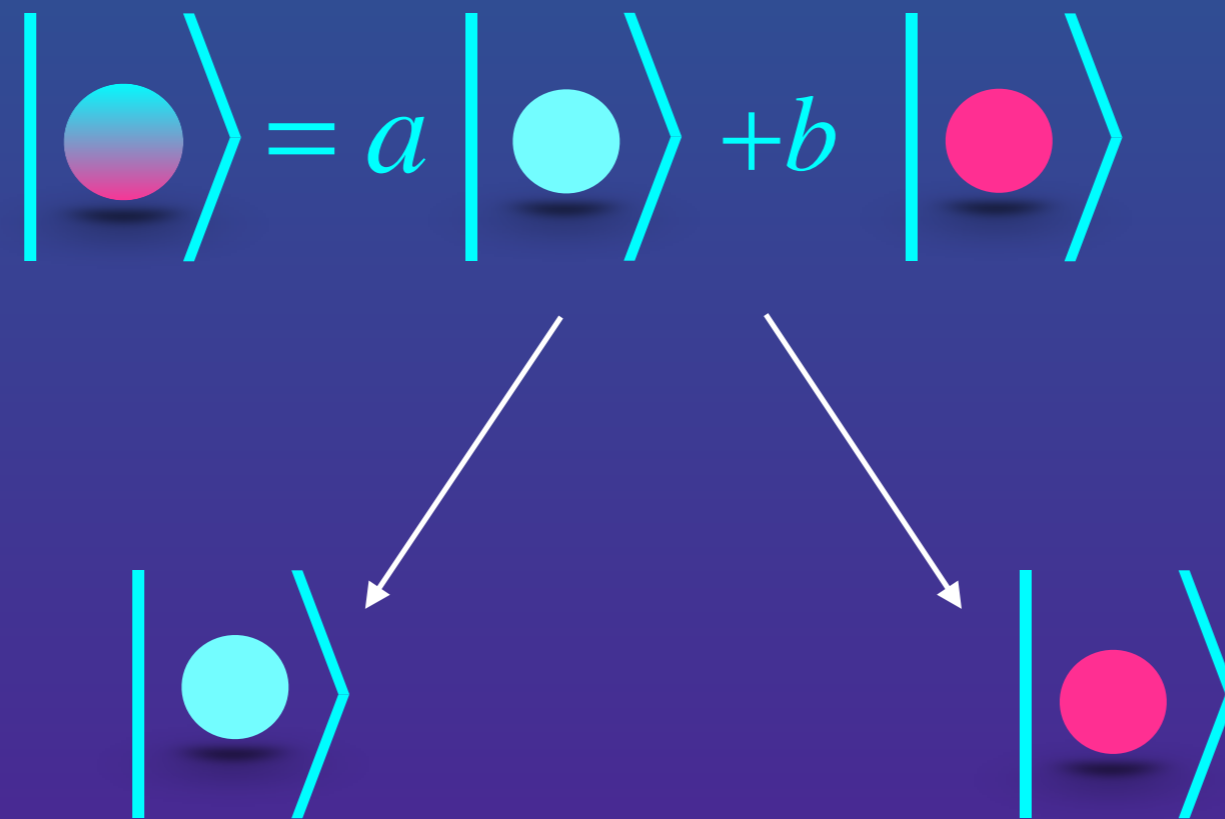


$$|\psi\rangle \otimes |0\rangle \rightarrow |\psi\rangle \otimes |\psi\rangle$$

Quantum Errors are continuous

$$|\text{blue}\rangle = a |\text{cyan}\rangle + b |\text{magenta}\rangle \longrightarrow |\text{blue}\rangle = a' |\text{cyan}\rangle + b' |\text{magenta}\rangle$$

They cannot be observed

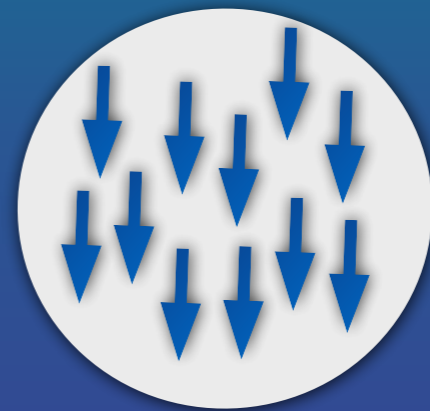


Topological Qubits

Merging the good features of both



$|\bar{0}\rangle$



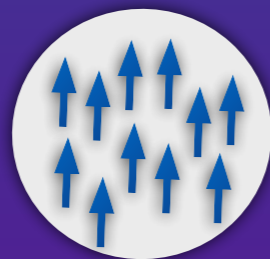
$|\bar{1}\rangle$

$$a|\bar{0}\rangle + b|\bar{1}\rangle$$

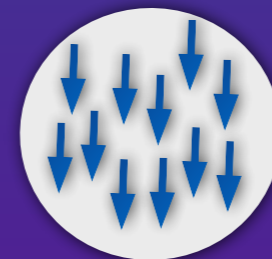
Ising Model

$$H = - \sum_{\langle i, j \rangle} z_i z_j$$

$$|\bar{0}\rangle = |\uparrow\uparrow\uparrow \dots \uparrow\uparrow\uparrow\rangle$$

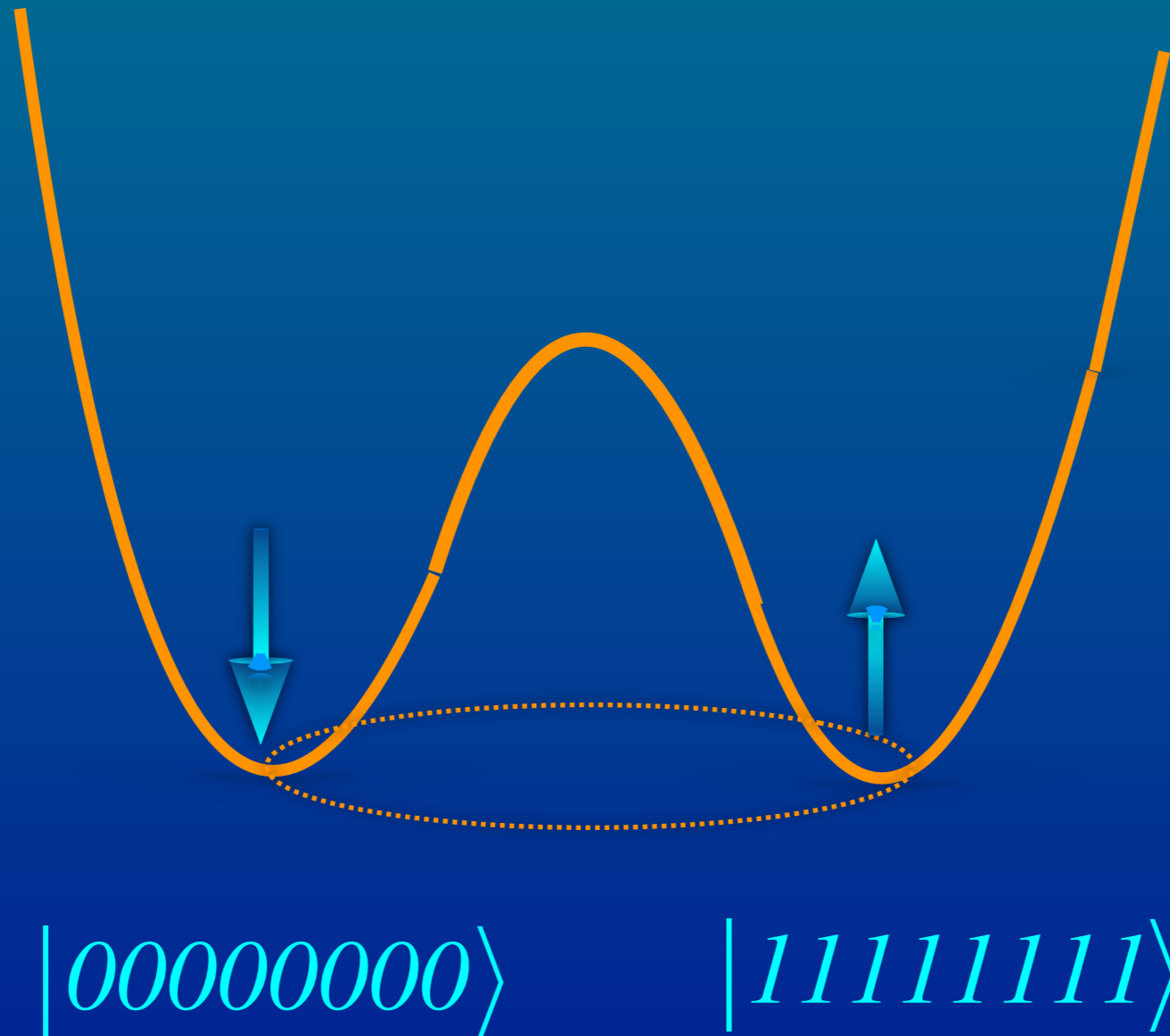


$$|\bar{1}\rangle = |\downarrow\downarrow\downarrow \dots \downarrow\downarrow\downarrow\rangle$$



Symmetry Breaking

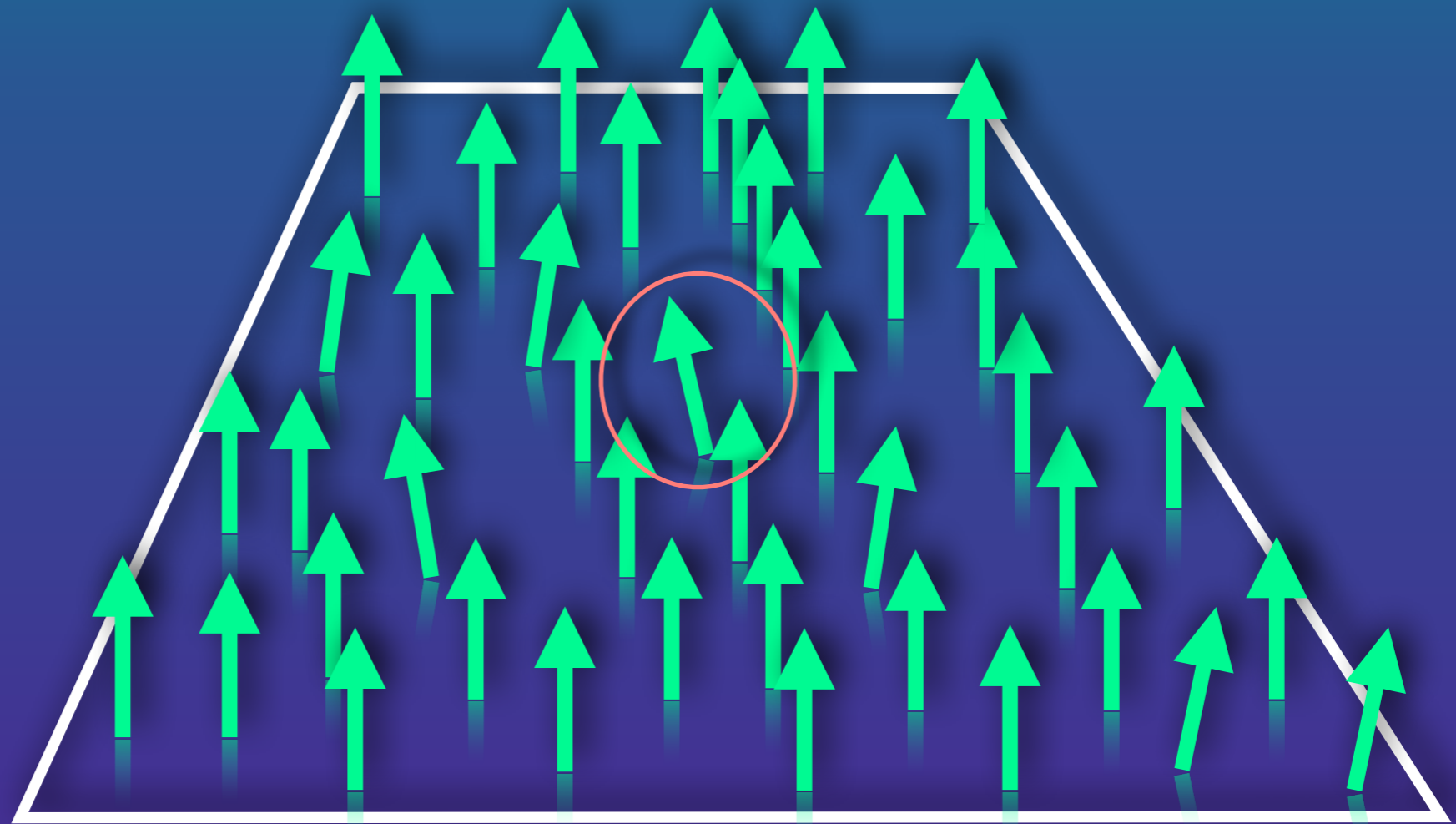
$$H = - \sum_{\langle i,j \rangle} z_i z_j$$



Local Order

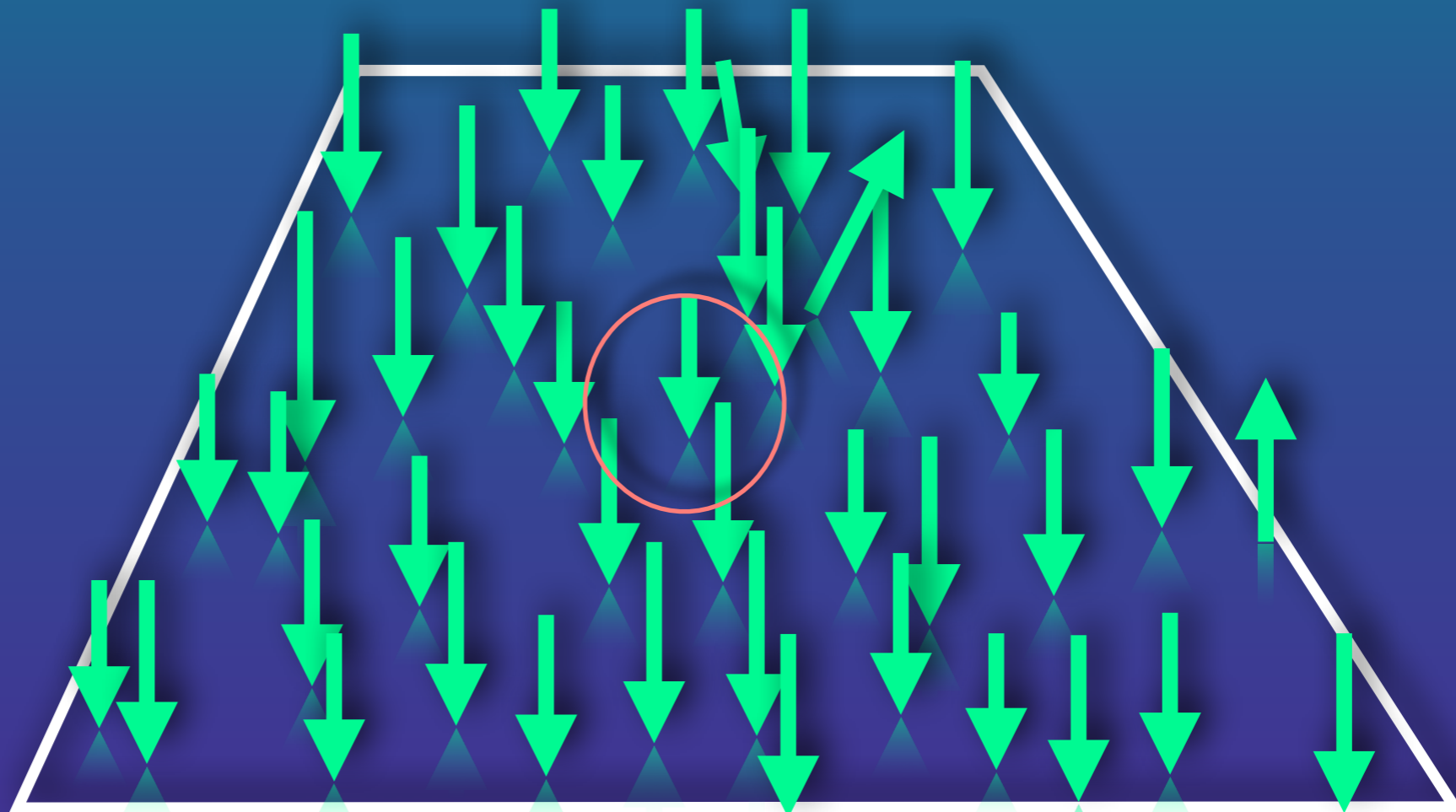
$$\langle \sigma_z \rangle = 1$$

$|0\rangle$



$$\langle \sigma_z \rangle = -1$$

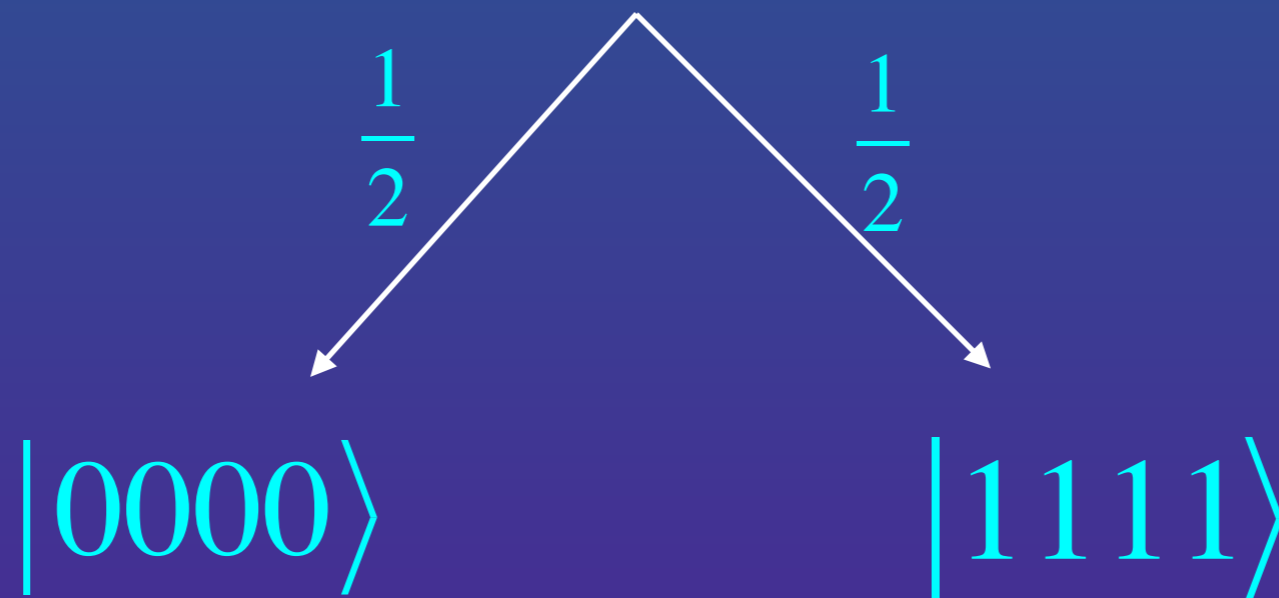
$|1\rangle$



But local order cannot
produce a topological
qubit!

Local order is extremely fragile

$$|GHZ\rangle = \frac{1}{\sqrt{2}}(|0000\rangle + |1111\rangle)$$



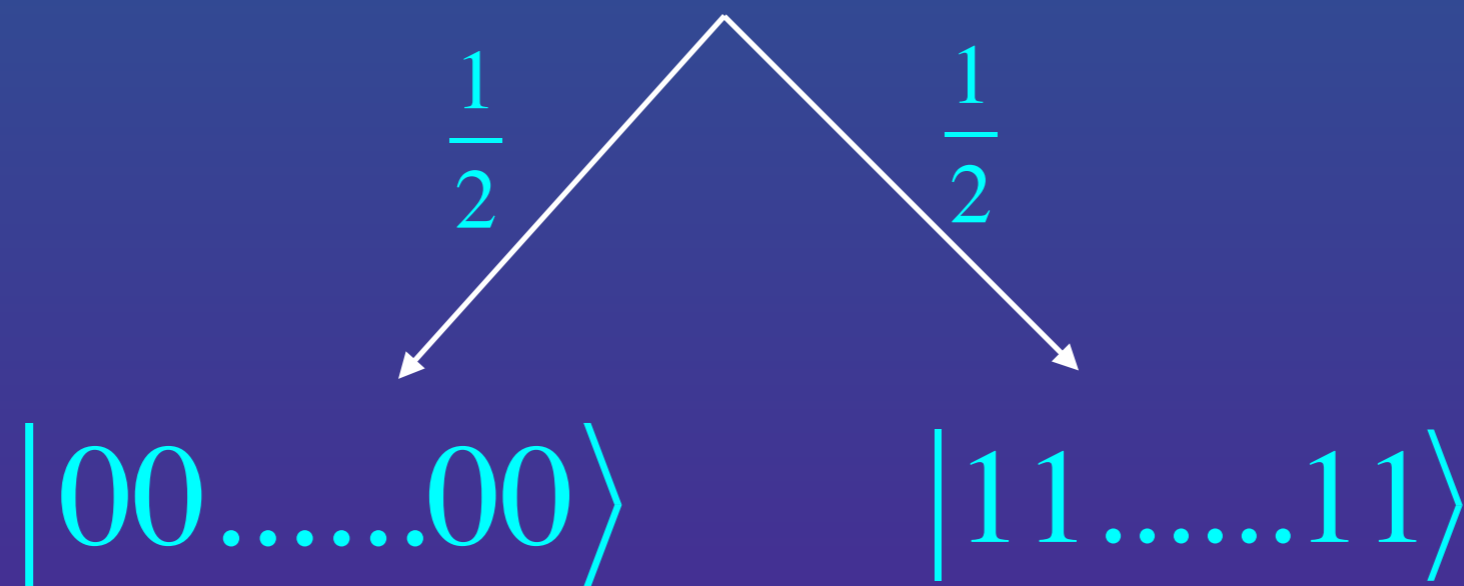
$$|W\rangle = \frac{1}{2}(|1000\rangle + |0100\rangle + |0010\rangle + |0001\rangle)$$



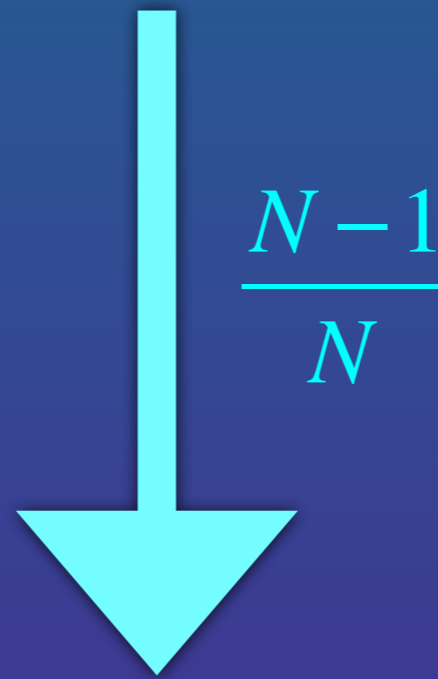
$$|W_1\rangle = |000\rangle$$

$$|W_0\rangle = \frac{1}{\sqrt{3}}(|100\rangle + |010\rangle + |001\rangle)$$

$$|GHZ\rangle = \frac{1}{\sqrt{2}} (|00\dots 00\rangle + |11\dots 11\rangle)$$



$$|W\rangle = \frac{1}{\sqrt{N}} (|100\dots000\rangle + |010\dots000\rangle + |001\dots000\rangle + \dots + |00\dots001\rangle)$$



$$|W_0\rangle = \frac{1}{\sqrt{N-1}} (|100\dots00\rangle + |010\dots00\rangle + |001\dots00\rangle + \dots + |00\dots01\rangle)$$

What we want?

$$|\text{blue}\rangle = a |\text{red}\rangle + b |\text{green}\rangle$$

- 1-Degenerate ground state
- 2-Existence of Gap
- 3-Not locally distinguishable
- 4-Robust to perturbations

A system with degenerate ground state



$|\psi_0\rangle$



$|\psi_1\rangle$

which cannot be distinguished,
by any local observable!

$$\langle \psi_0 | K | \psi_0 \rangle = \langle \psi_1 | K | \psi_1 \rangle$$

The Stabilizer Formalism

$$|GHZ\rangle = |000\rangle + |111\rangle$$

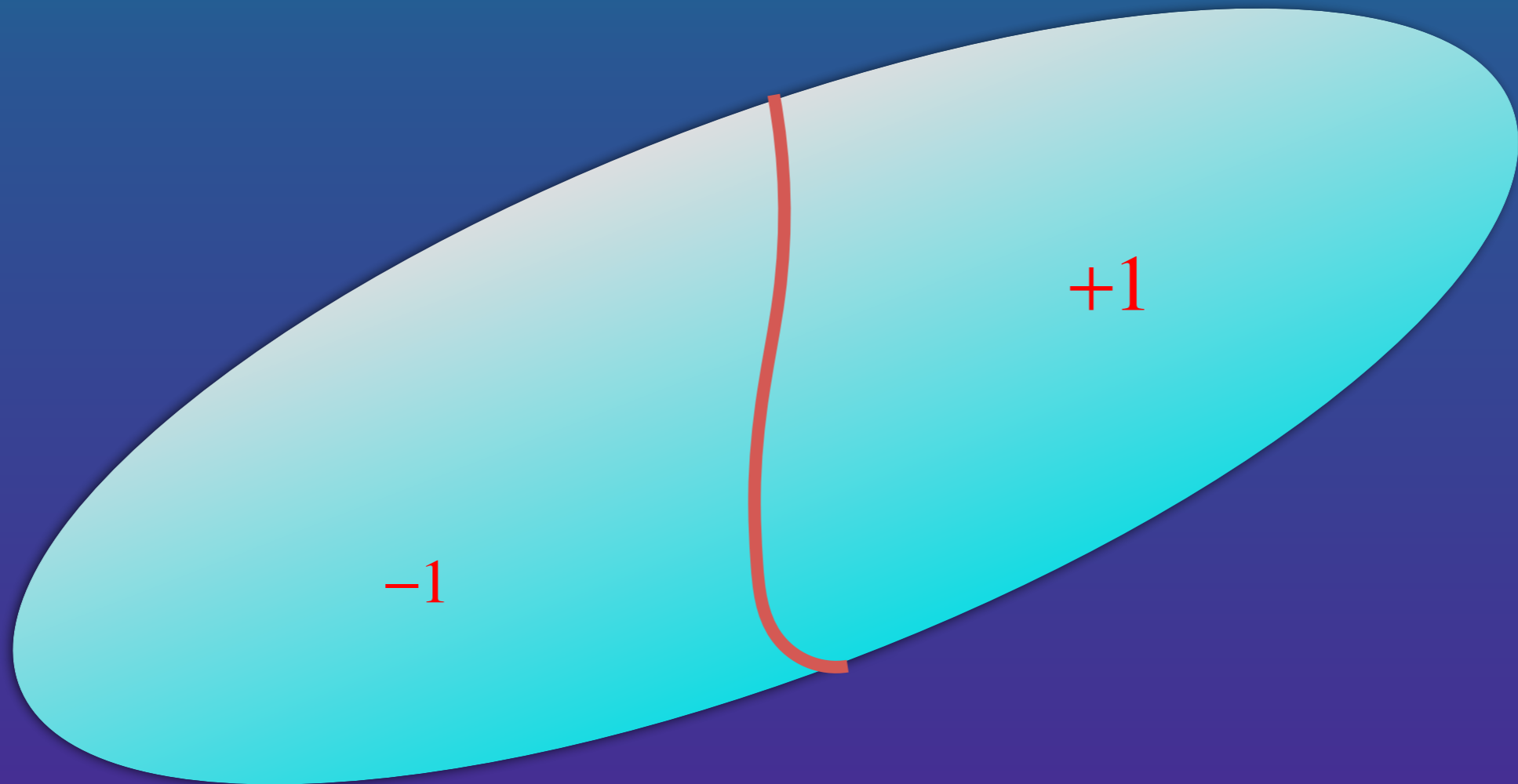
$$X_1 X_2 X_3 |GHZ\rangle = |GHZ\rangle$$

$$Z_1 Z_2 |GHZ\rangle = |GHZ\rangle$$

$$S = \{I, Z_1 Z_2, Z_1 Z_3, X_1 X_2 X_3\}$$

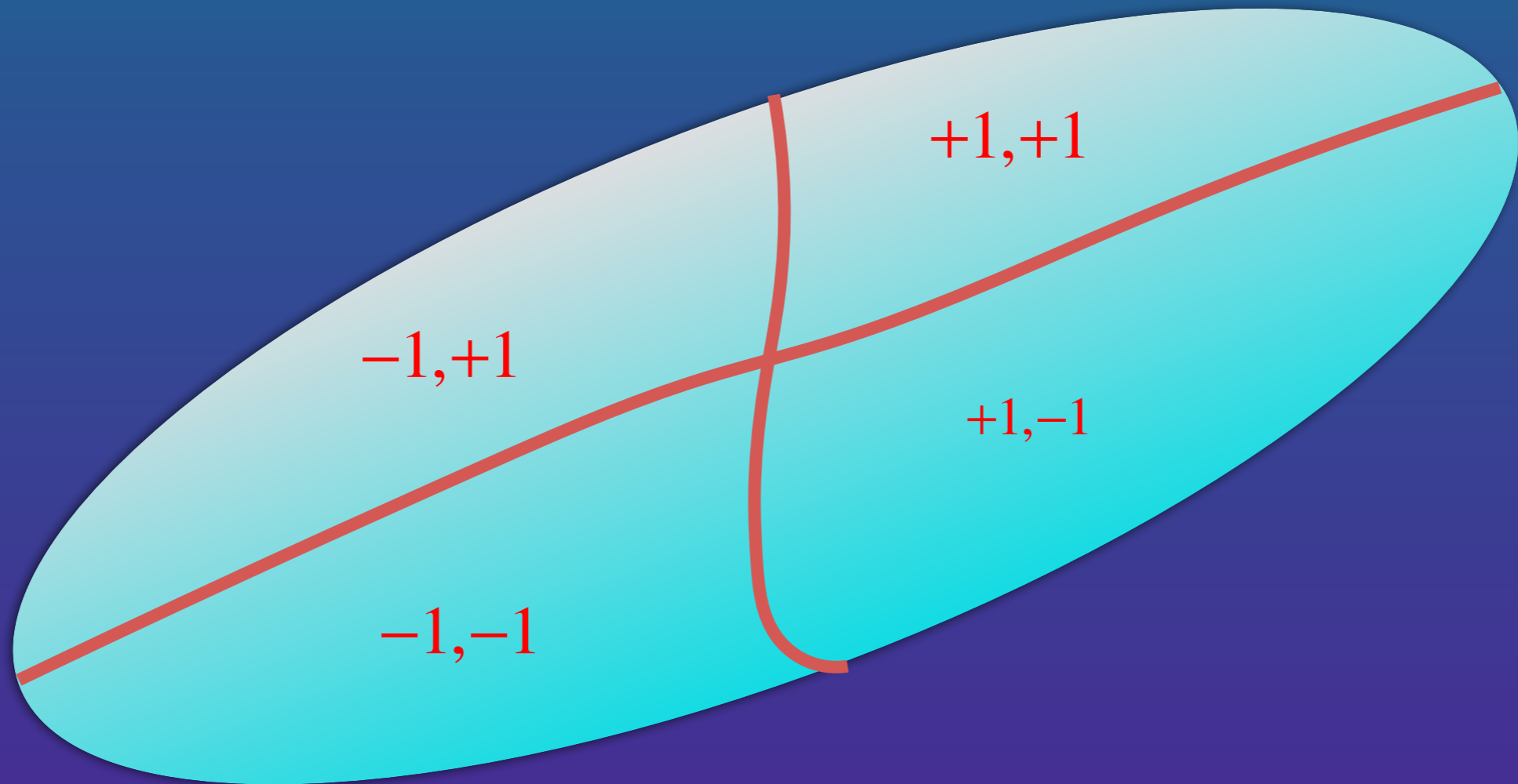
The Stabilizer Formalism

S_1



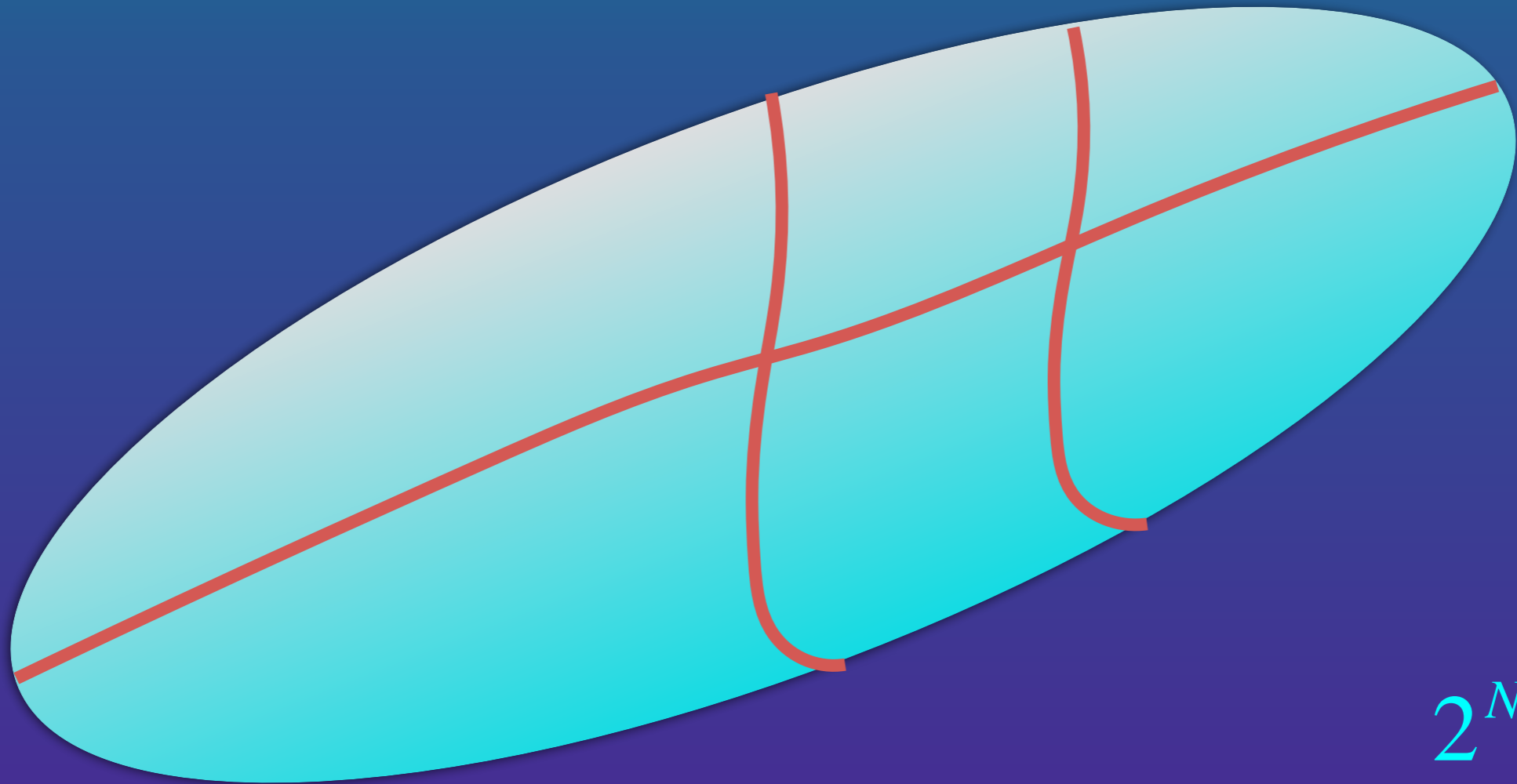
The Stabilizer Formalism

S_1 S_2



The Stabilizer Formalism

S_1 S_2 S_3

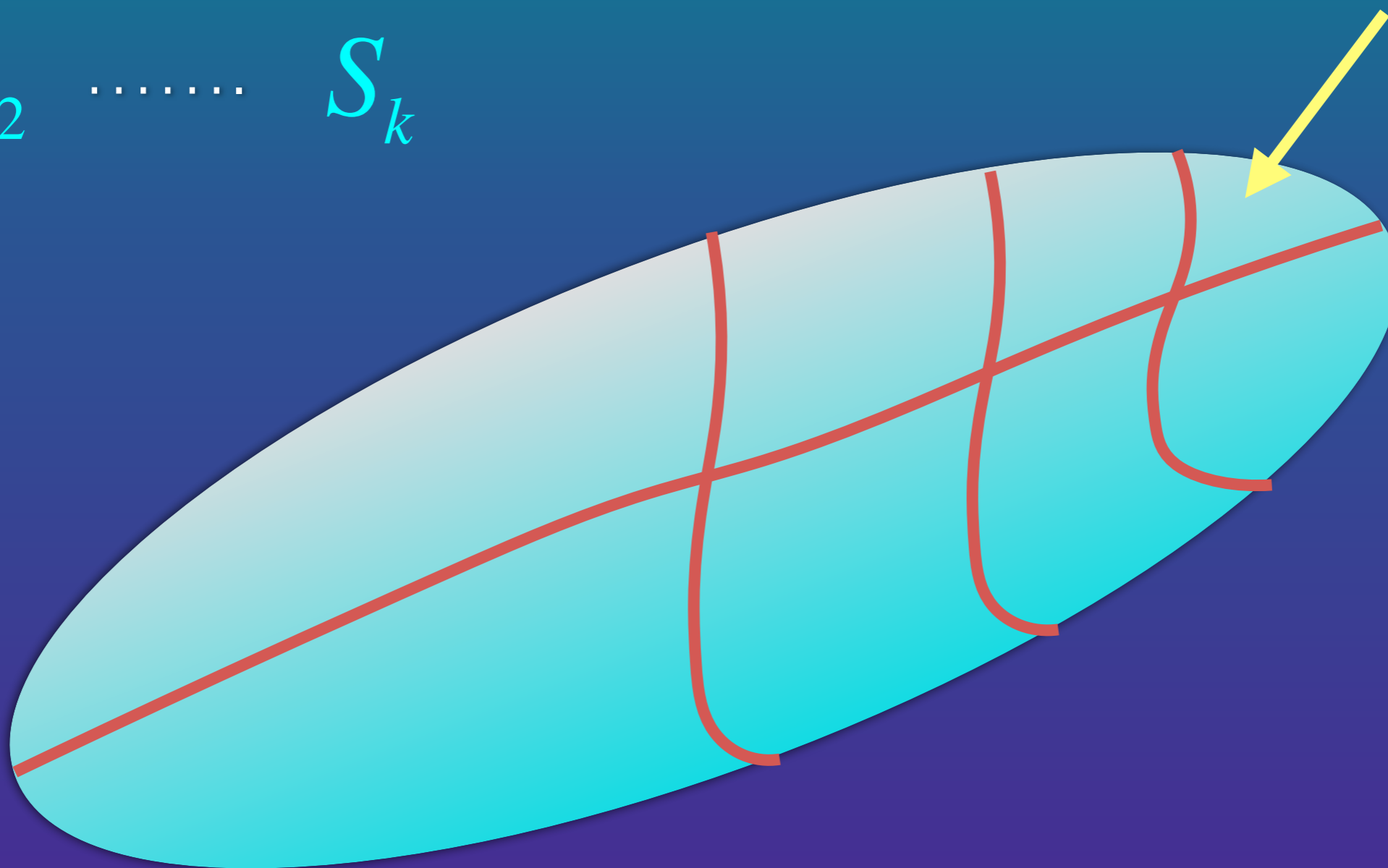


$$\frac{2^N}{2 * 2 * 2}$$

The Stabilizer Formalism

S_1 S_2 S_k

$$\frac{2^N}{2^k}$$



The Hamiltonian

$$H = -\mathcal{S}_1 - \mathcal{S}_2 - \dots - \mathcal{S}_k$$

All the local operators



$$H|\psi\rangle = -k|\psi\rangle$$

A ground state



The order of degeneracy

$$\frac{2^N}{2^k} = 2^{N-k}$$

The Kitaev Model



Notations

$$x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{cases} |+\rangle \\ |-\rangle \end{cases}$$

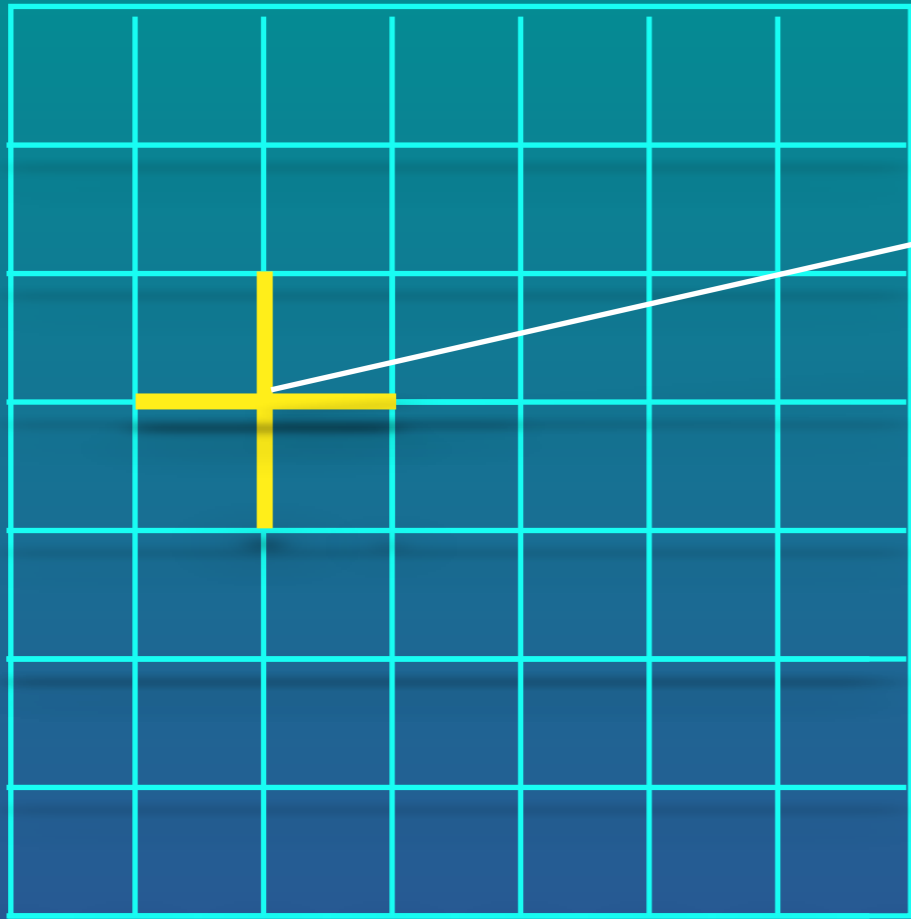
$$z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{cases} |0\rangle \\ |1\rangle \end{cases}$$

$$z|+\rangle = |-\rangle$$

$$x|0\rangle = |1\rangle$$

$$z|-\rangle = |+\rangle$$

$$x|1\rangle = |0\rangle$$



$$A_s = x_1 x_2 x_3 x_4$$

$$A_s^2 = I$$

$$\prod_s A_s = I$$

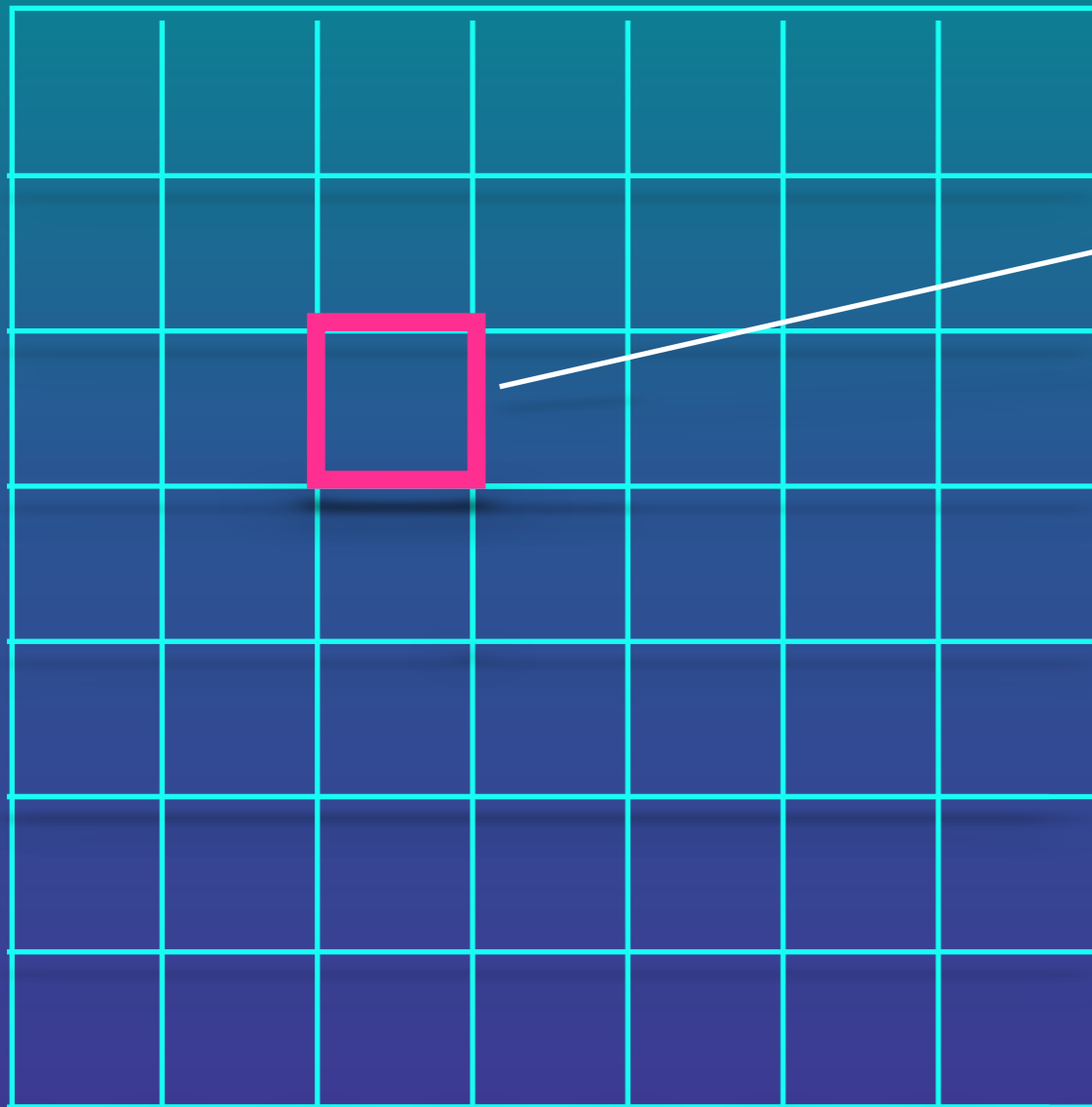
Number of vertices=N

Number of links=2N

Dimension of Hilbert Space = 2^{2N}

Number of independent A's = N-1

$$\text{Degeneracy} = \frac{2^{2N}}{2^{N-1}} = 2^{N+1}$$



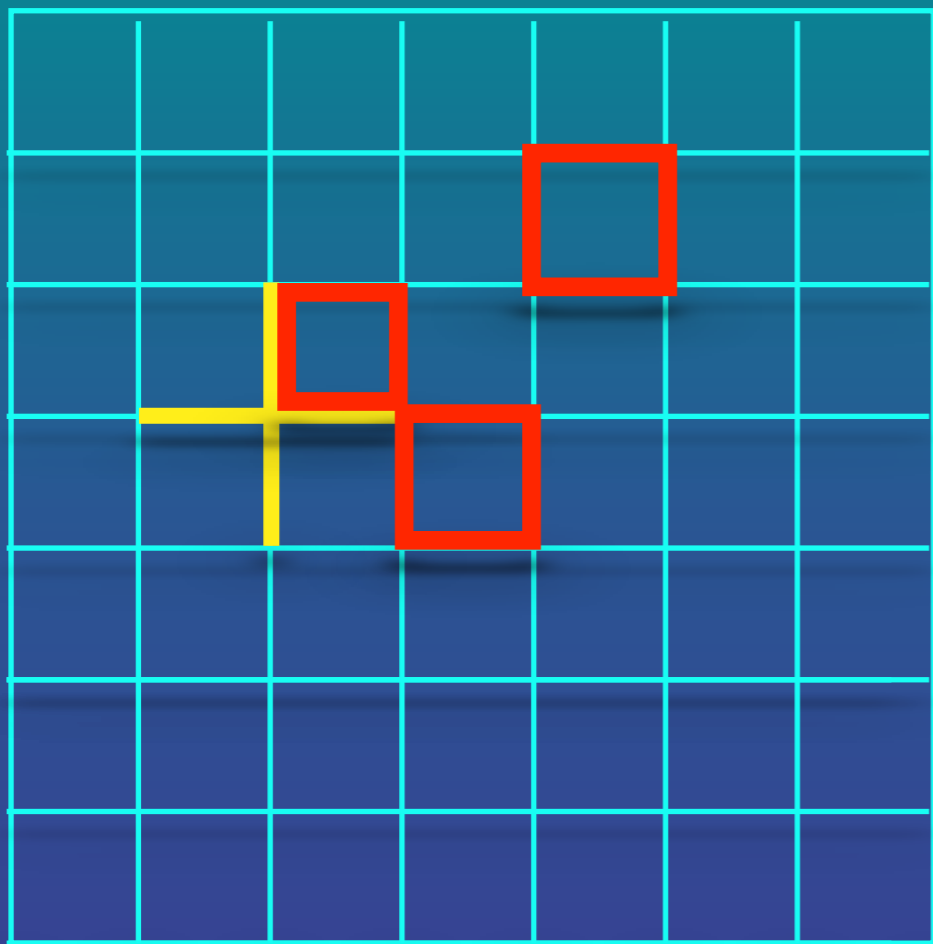
$$B_p = z_1 z_2 z_3 z_4$$

$$B_p^2 = I$$

$$\prod_p B_p = I$$

Number of faces = N

Number of Independent B's = $N-1$



$$[A_s, B_p] = 0$$

$$H = -\sum_s A_s - \sum_p B_p$$

$$\text{Degeneracy} = \frac{2^{2N}}{2^{2N-2}} = 4$$

The ground state

$$H = -\sum_s A_s - \sum_p B_p$$

$$A_s |\phi\rangle = |\phi\rangle$$

$$B_p |\phi\rangle = |\phi\rangle$$

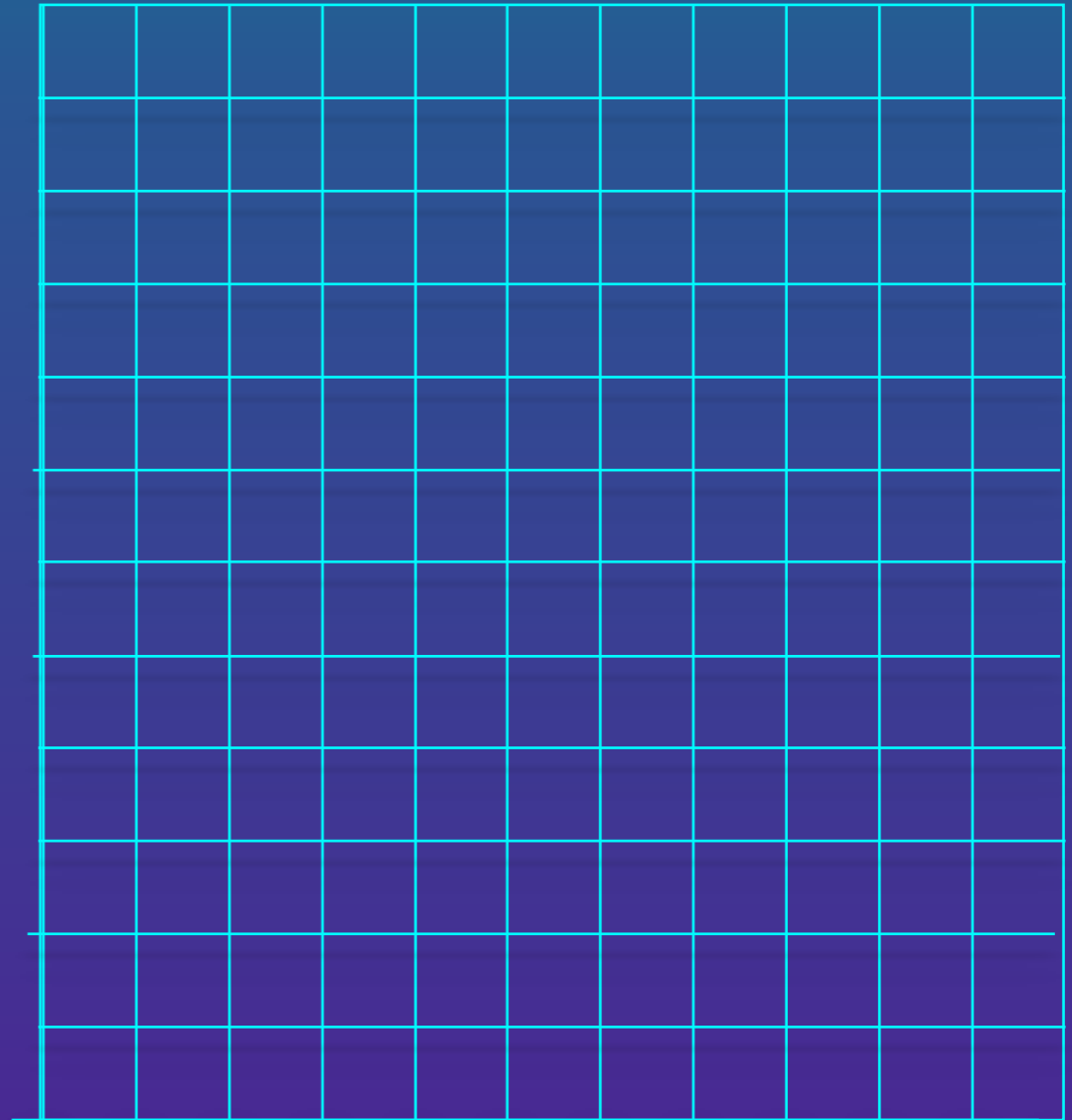
The ground state?

$$|\Omega\rangle = |+\rangle^{\otimes N}$$

$$A_s |\Omega\rangle = |\Omega\rangle$$

$$z|+\rangle = |-\rangle$$

$$B_p |\Omega\rangle \neq |\Omega\rangle$$



$$B_p(1 + B_p) = B_p + 1$$

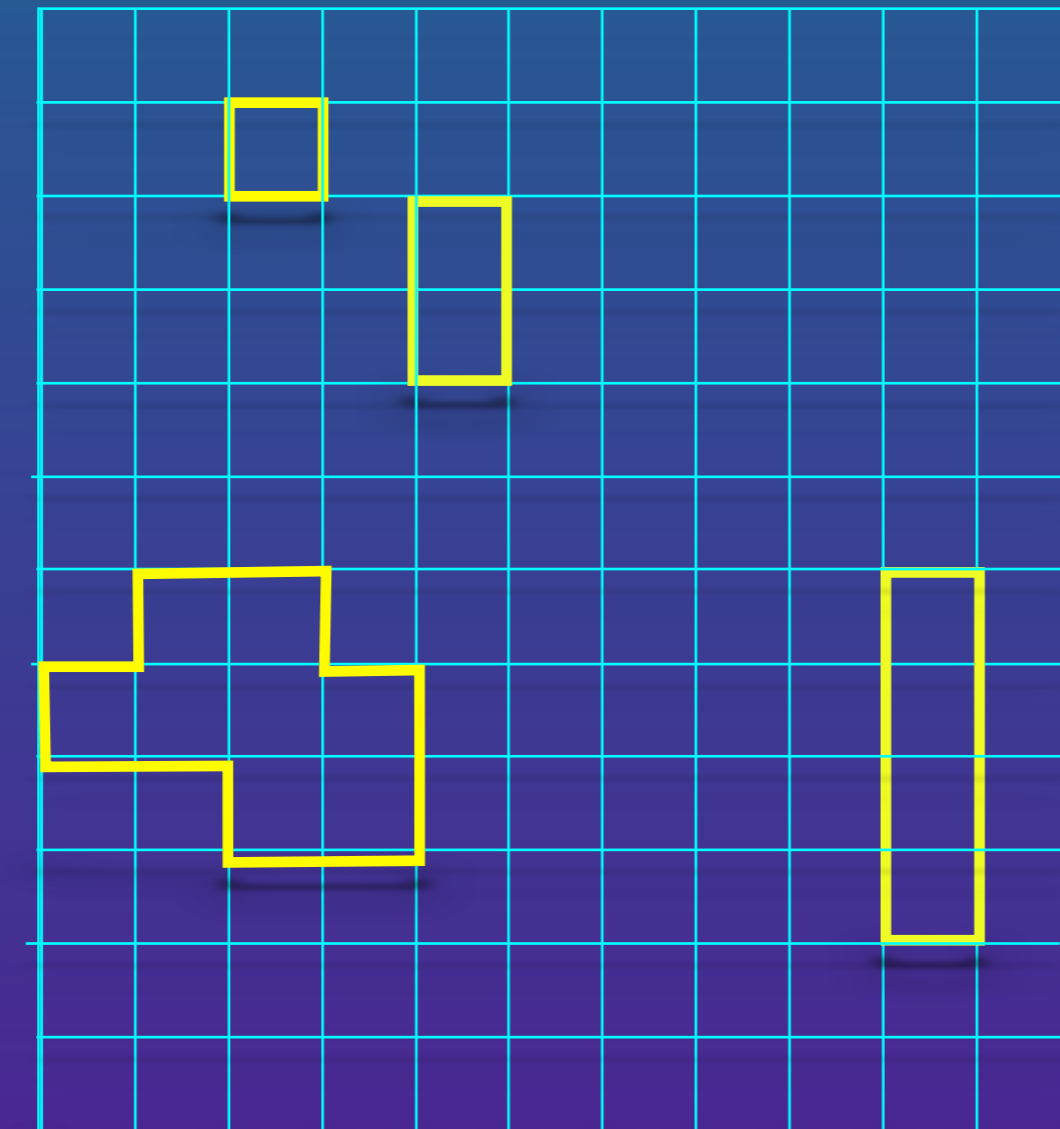
$$|\phi\rangle = (1 + B_p)|\Omega\rangle$$

$$A_s|\phi\rangle = |\phi\rangle$$

$$B_p|\phi\rangle = |\phi\rangle$$

$$|\varphi_0\rangle = \prod_p (1 + B_p)|\Omega\rangle$$

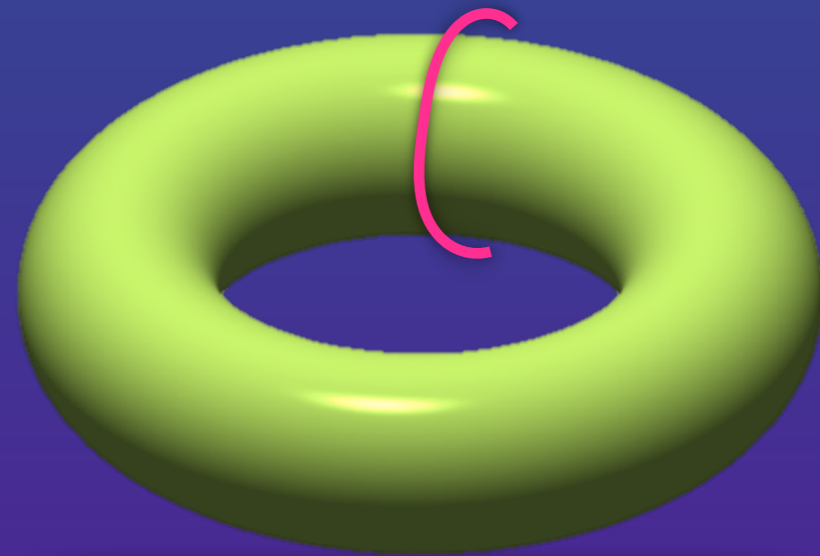
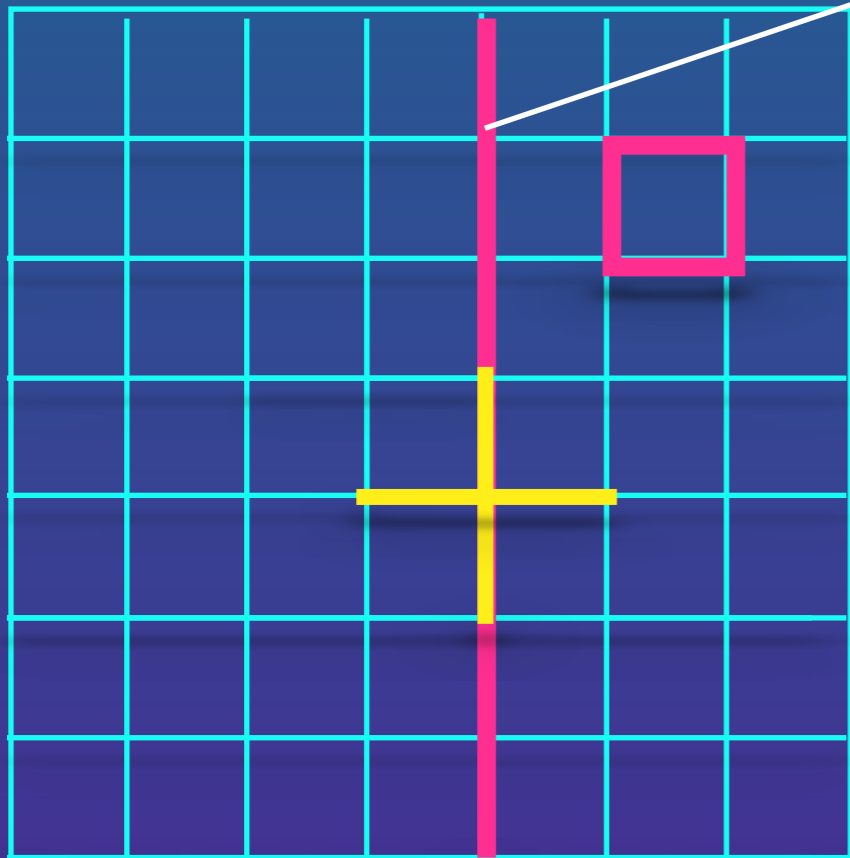
What the ground state looks like?



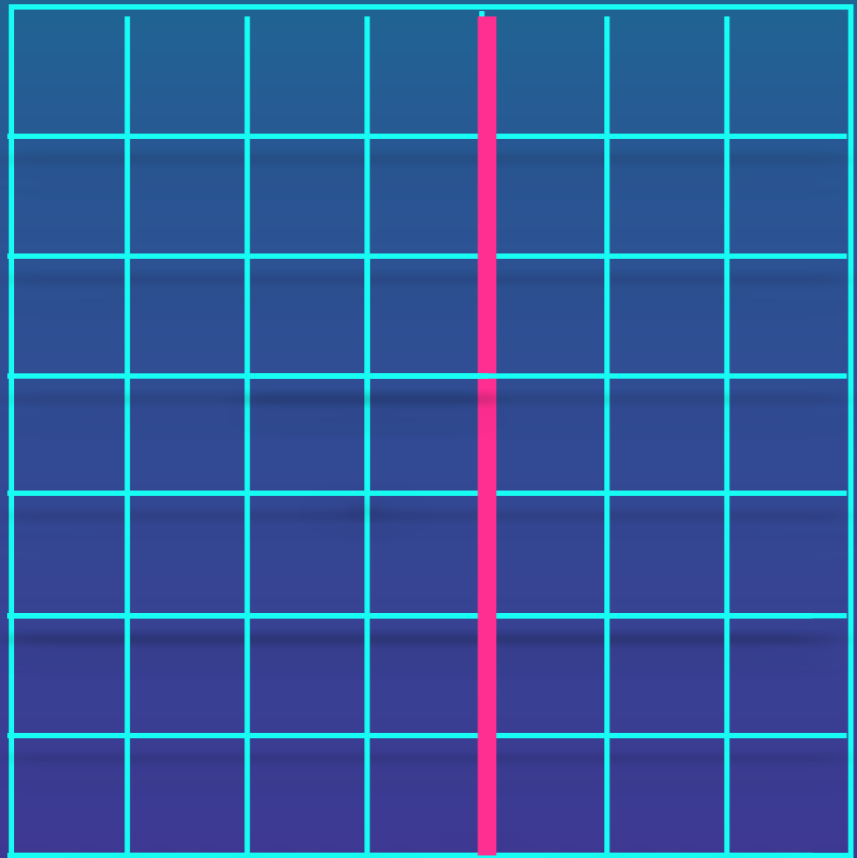
Where is the degeneracy?

$$Z_1 = \prod_{i \in C_1} z_i$$

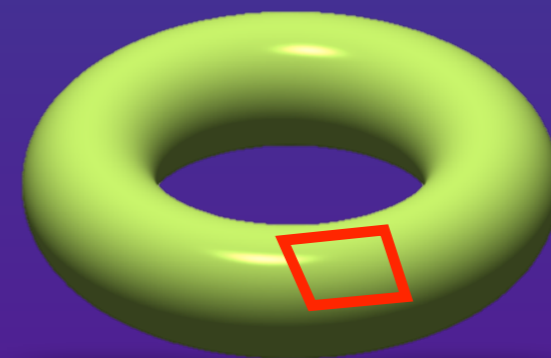
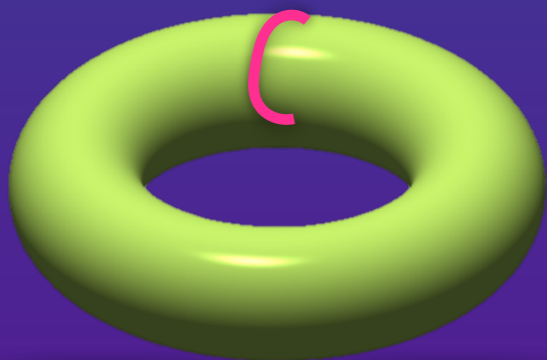
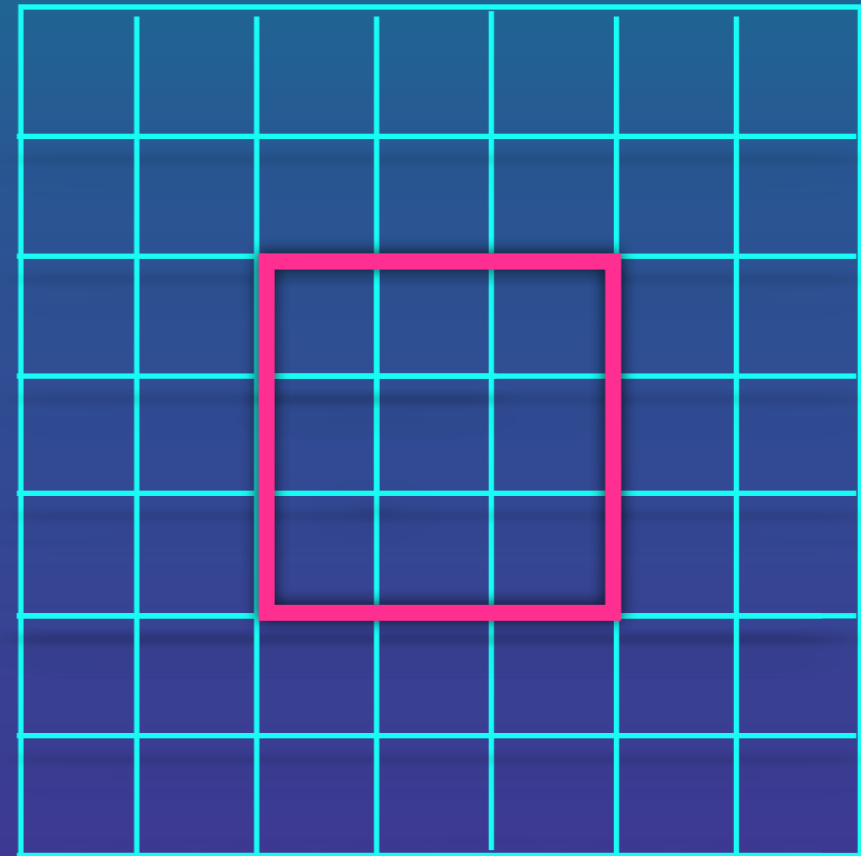
$$[Z_1, H] = 0$$

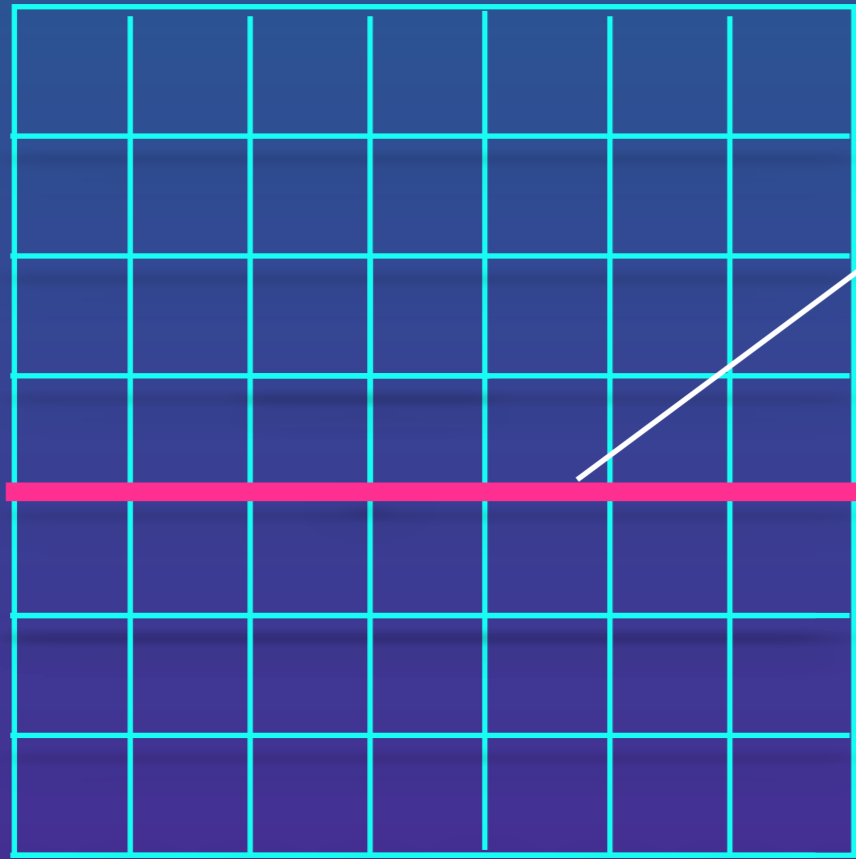


Where is the degeneracy?



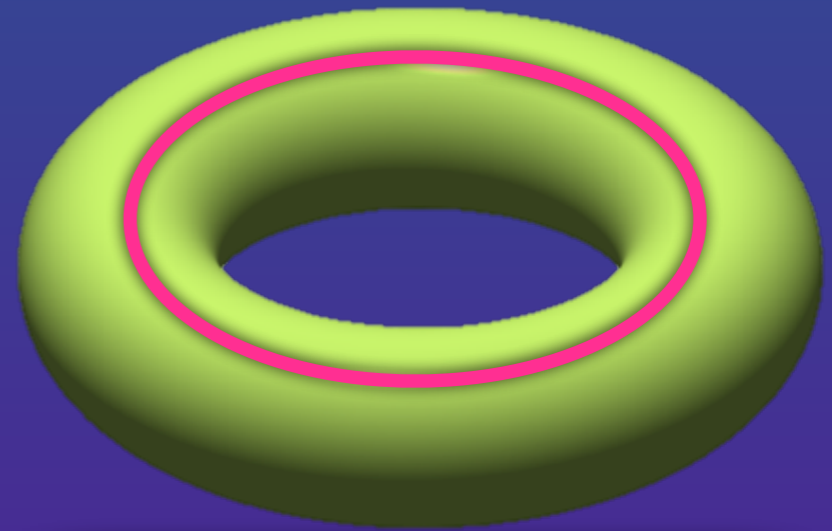
\neq



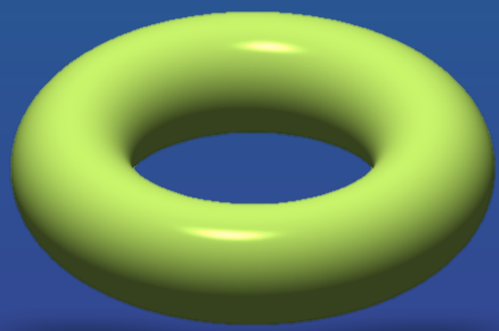


$$Z_2 = \prod_{i \in C_2} z_i$$

$$[Z_2, H] = 0$$



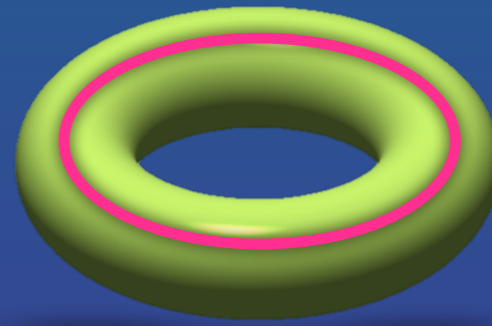
Four ground states



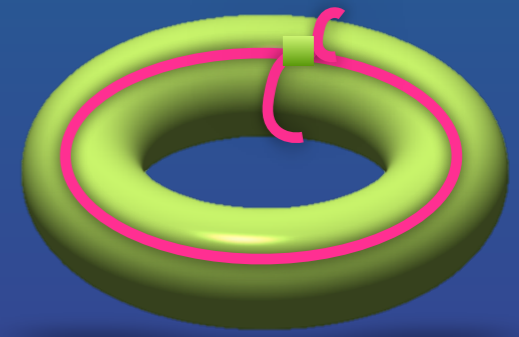
$$|\phi_{00}\rangle$$



$$|\phi_{01}\rangle = Z_1 |\phi_{00}\rangle$$



$$|\phi_{10}\rangle = Z_2 |\phi_{00}\rangle$$



$$|\phi_{11}\rangle = Z_1 Z_2 |\phi_{00}\rangle$$

We achieved goal no. 1

Degenerate Ground States.



We also achieved goal no. 2

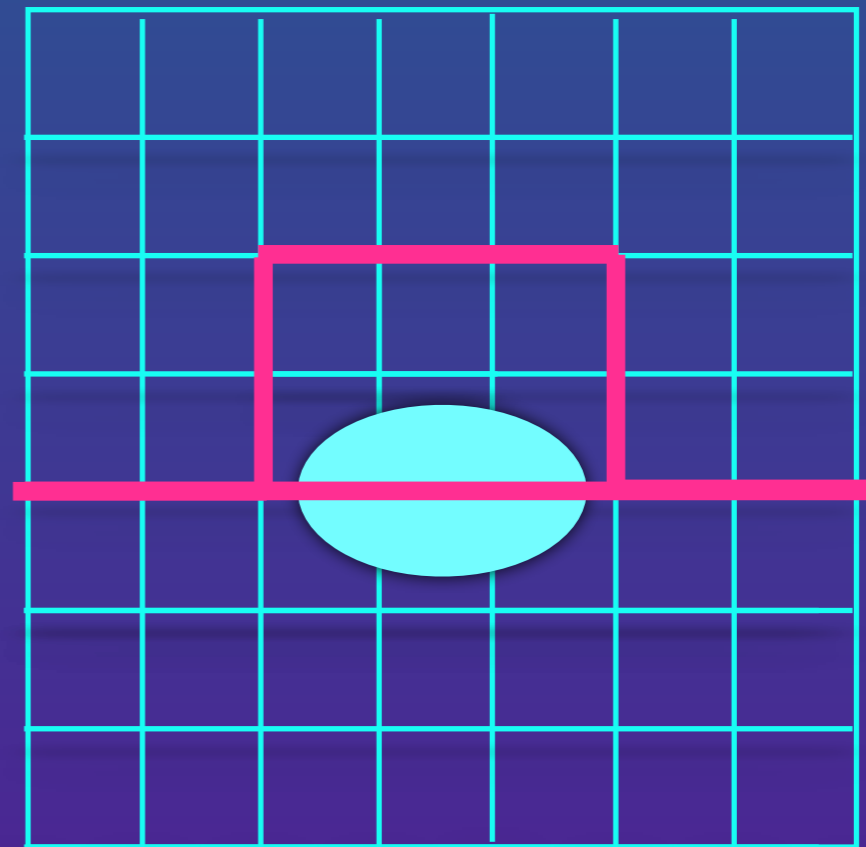
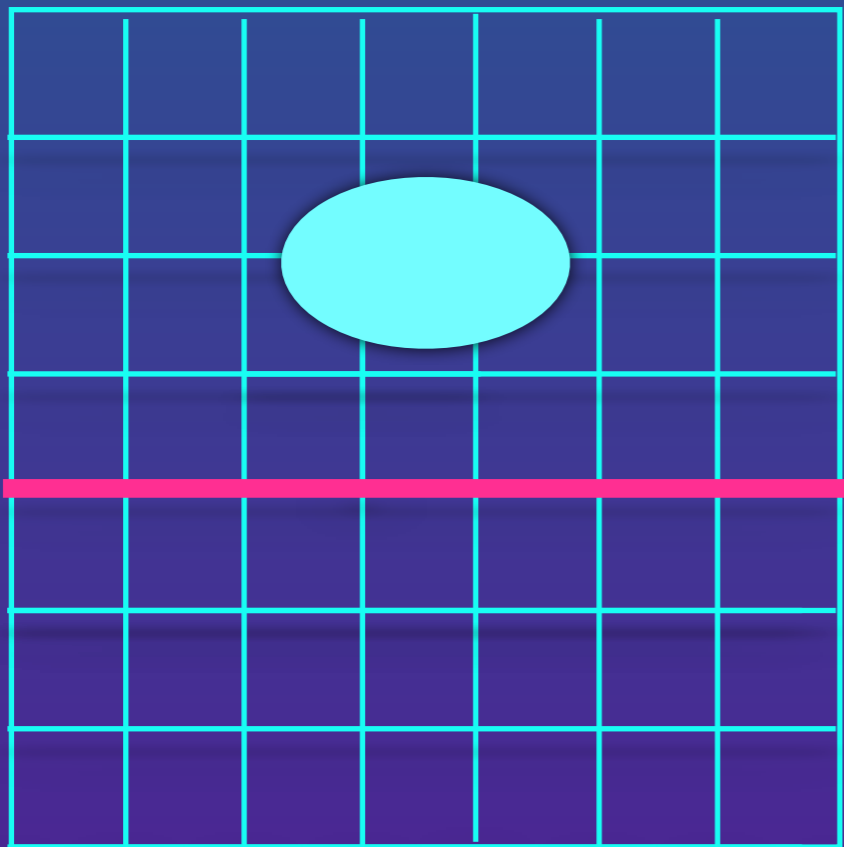
There is a finite Gap

$$H = -\sum_s A_s - \sum_p B_p$$



Can we distinguish the ground states
by local observation?

$$Z_1 O = O Z_1$$



Can we distinguish the ground states
by local observation?

$$\langle \phi | O | \phi \rangle = \langle \phi' | O | \phi' \rangle$$

$$\langle \phi' | O | \phi' \rangle = \langle \phi | Z_1 O Z_1 | \phi \rangle = \langle \phi | Z_1 Z_1 O | \phi \rangle = \langle \phi | O | \phi \rangle$$

So we have also achieved goal 3:

The states are not distinguishable locally.

What about goal 4?

What happens if I perturb the Hamiltonian?

$$H \rightarrow H + \sum_i O_i$$

$$\Delta E_\alpha = \langle \psi_\alpha | \sum_i O_i | \psi_\alpha \rangle$$



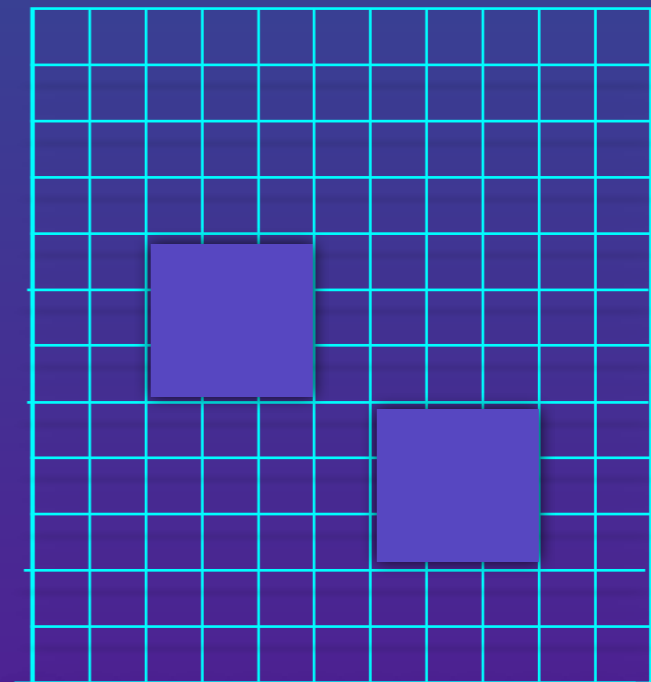
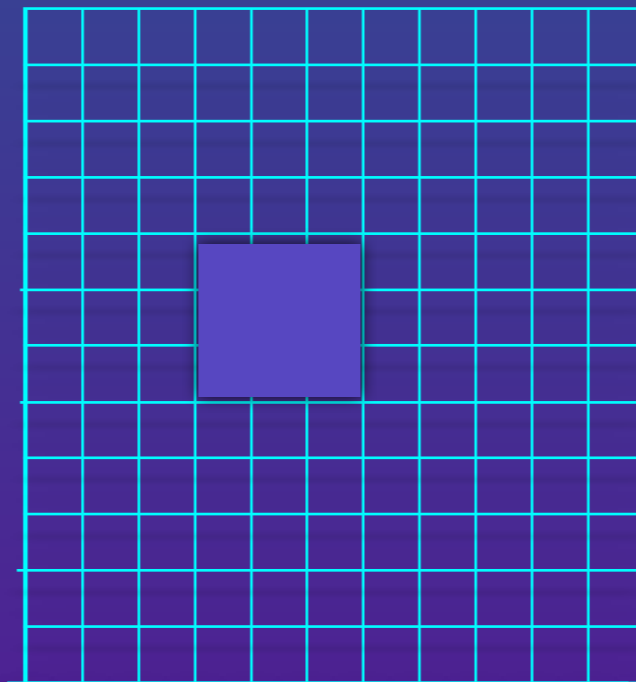
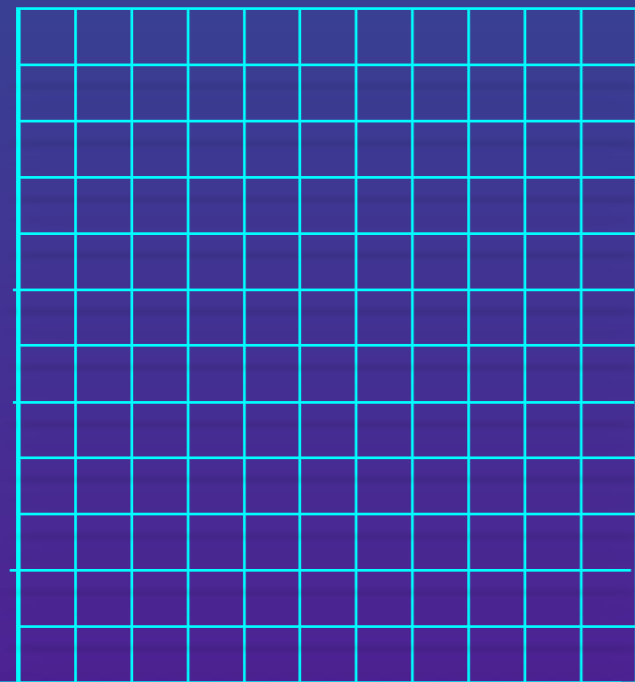
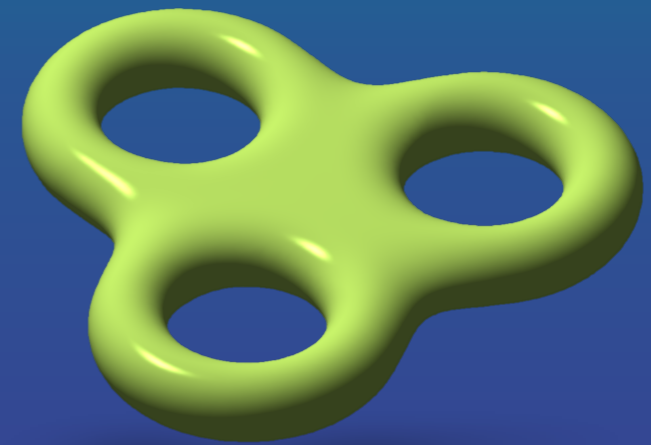
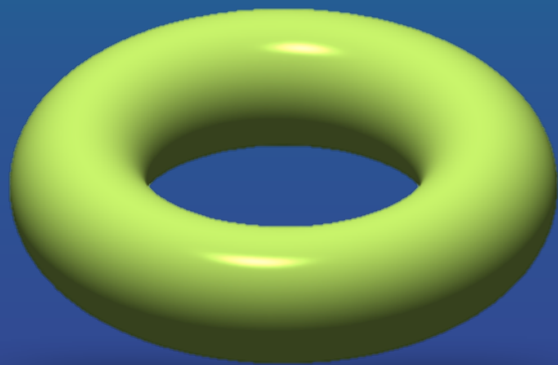
$$\Delta E_\alpha = \Delta E_\beta$$

So we have also achieved goal 4:

The states are robust under perturbations.

We have two topological qubits.

More Qubits



What we have left out:

Anyonic Excitations

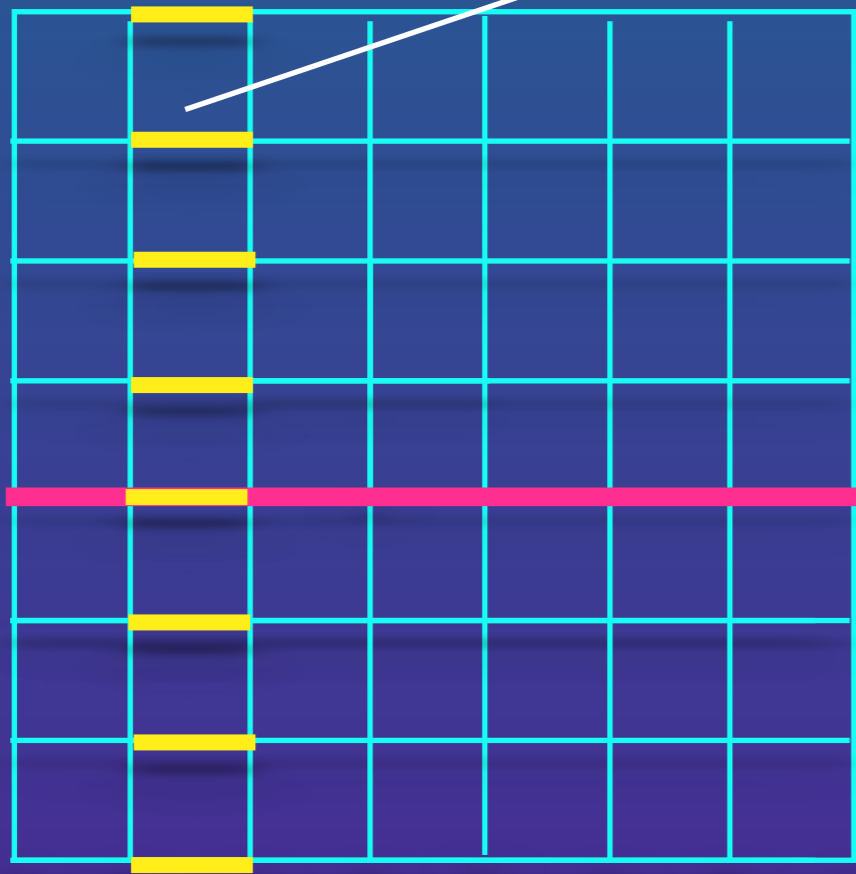
Moving anyons and performing gates

Non-Abelian Anyons and Universal quantum
computation

Thank you for your attention

String operators which distinguish the states!

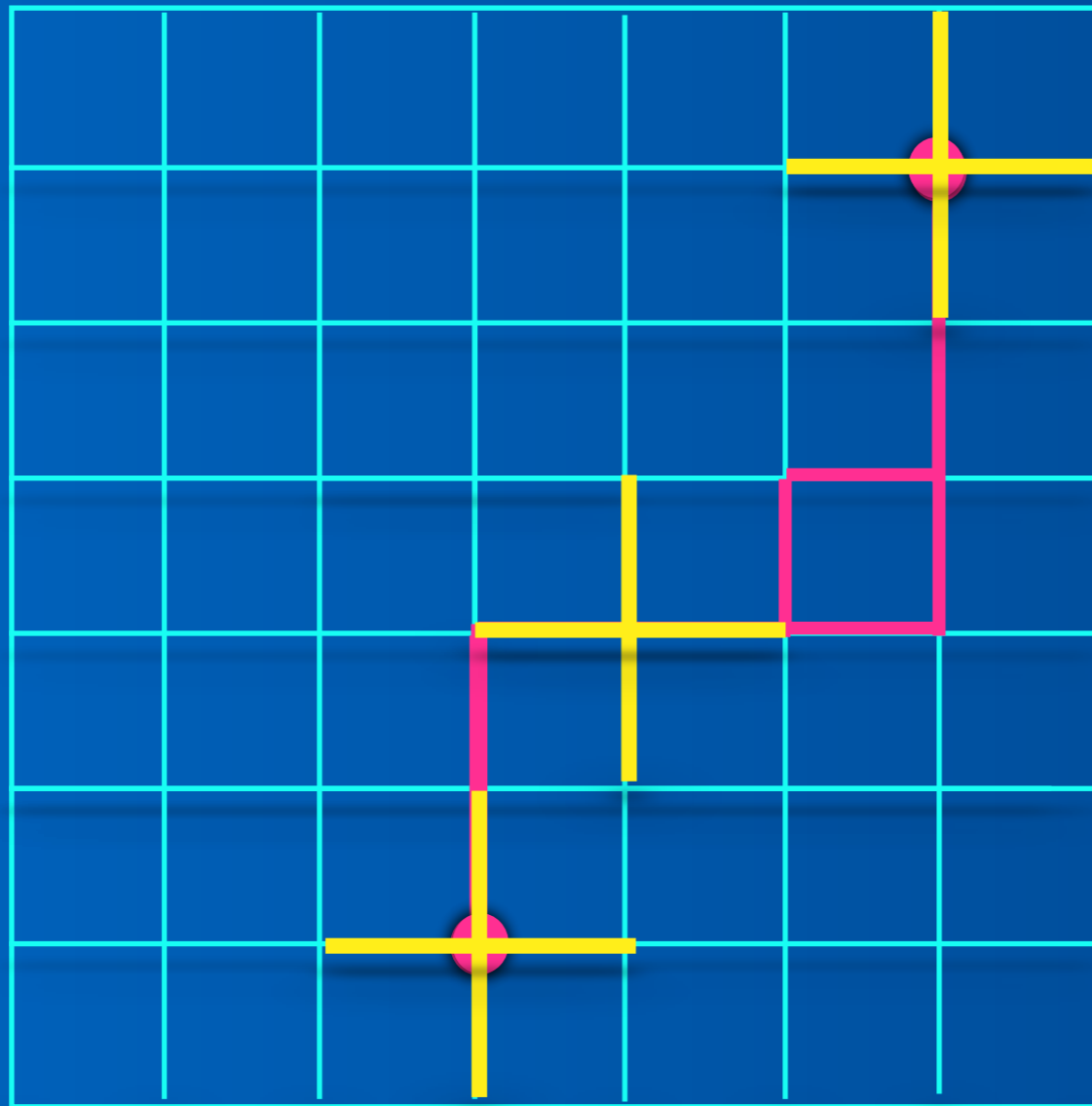
$$X_1 = \prod_{i=1}^N x_i$$



$$X_1 Z_1 = -Z_1 X_1$$

$$[X_1, H] = 0$$

Excited States: 1- Electric Anyons.

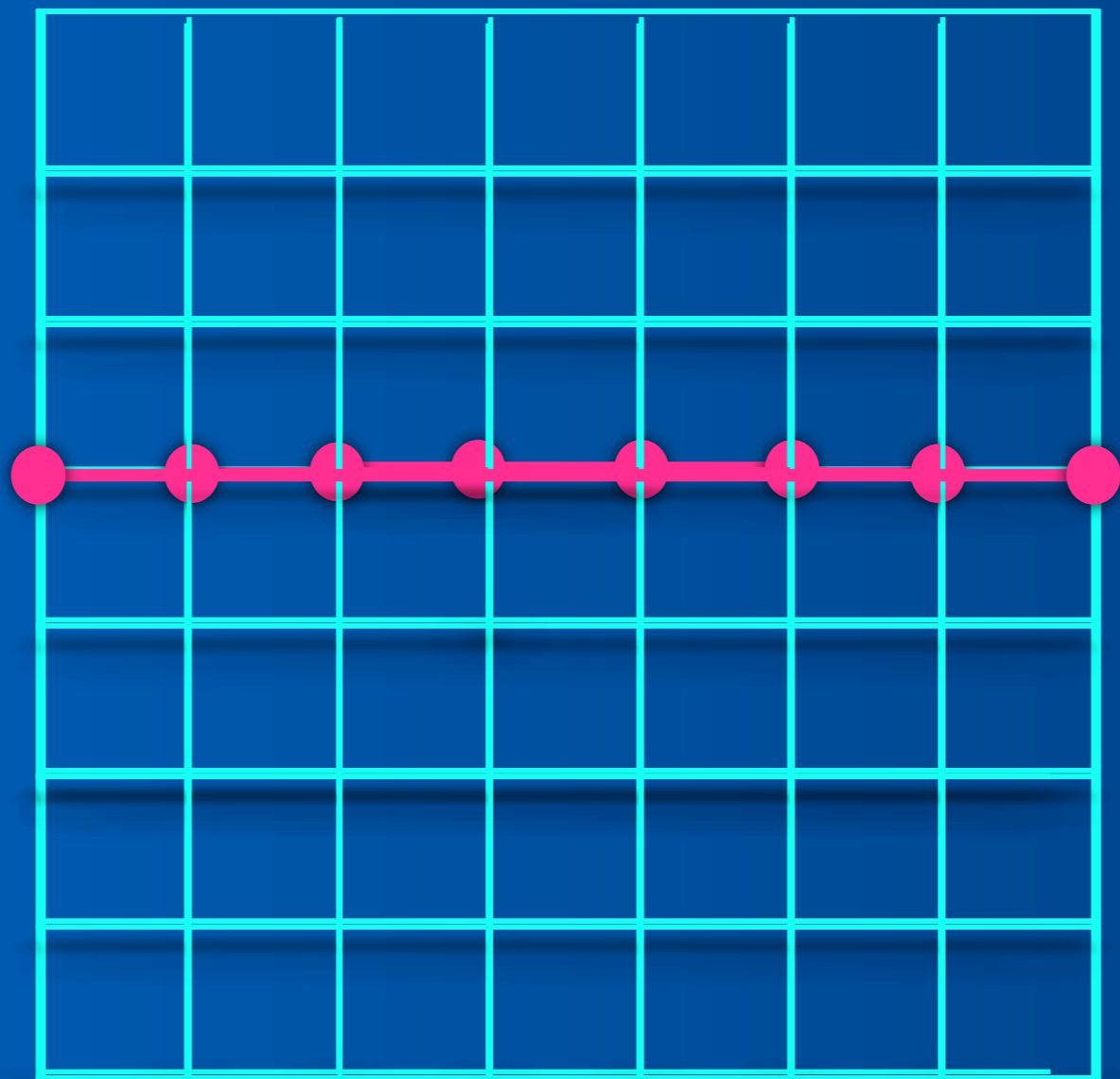


Each Anyon has an energy of
2 units.

Anyons are created in pairs.

The energy of the pair doesn't
depend on the path connecting them.

Another interpretation of String Operators

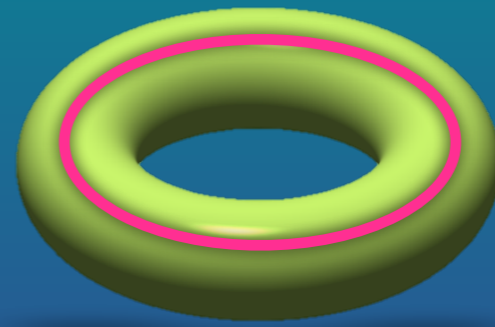


So by creating two electric Anyons,

Moving them across the Torus,

And annihilating them in the end,

We can implement a X gate on either of the qubits.



$$|\psi_{00}\rangle$$

$$|\psi_{10}\rangle = Z_1 |\psi_{00}\rangle$$

$$|\psi_{01}\rangle = Z_2 |\psi_{00}\rangle$$

$$|\psi_{11}\rangle = Z_2 Z_1 |\psi_{00}\rangle$$

 X_1

1

-1

1

-1

 X_2

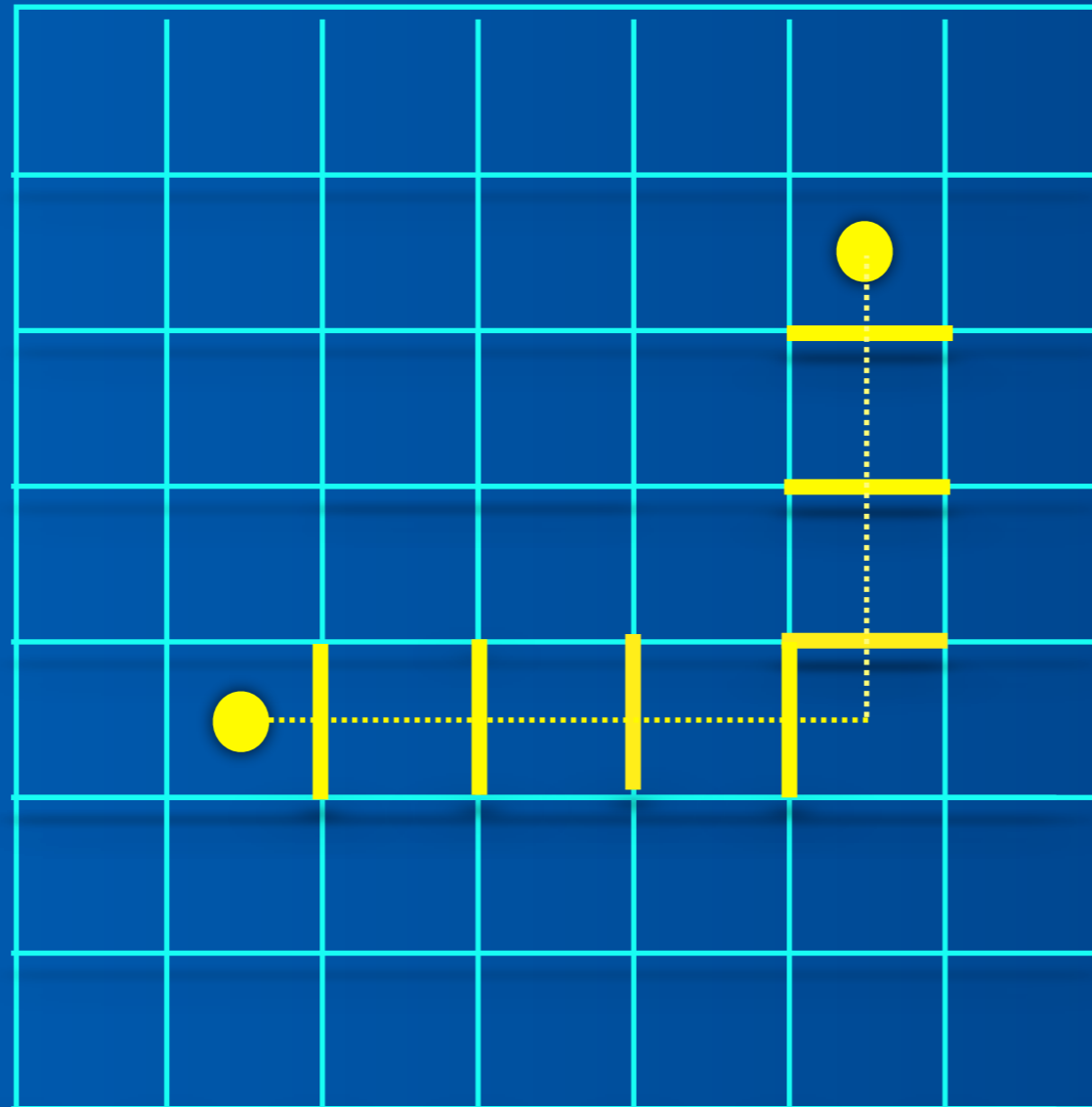
1

1

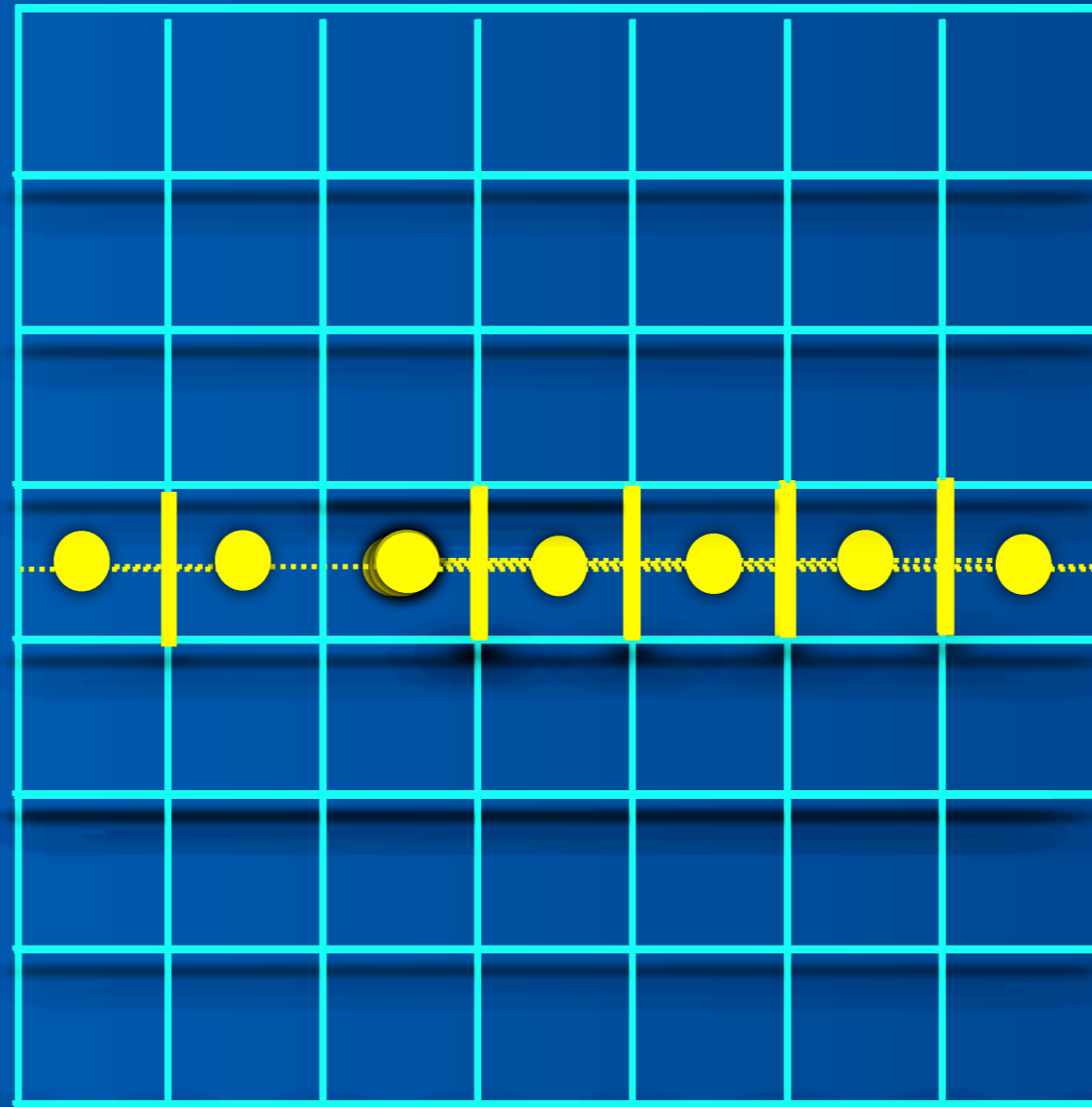
-1

-1

Excited States: Magnetic Anyons



Another interoperation of string operators



So by creating two Magnetic Anyons,

Moving them across the Torus,

And annihilating them in the end,

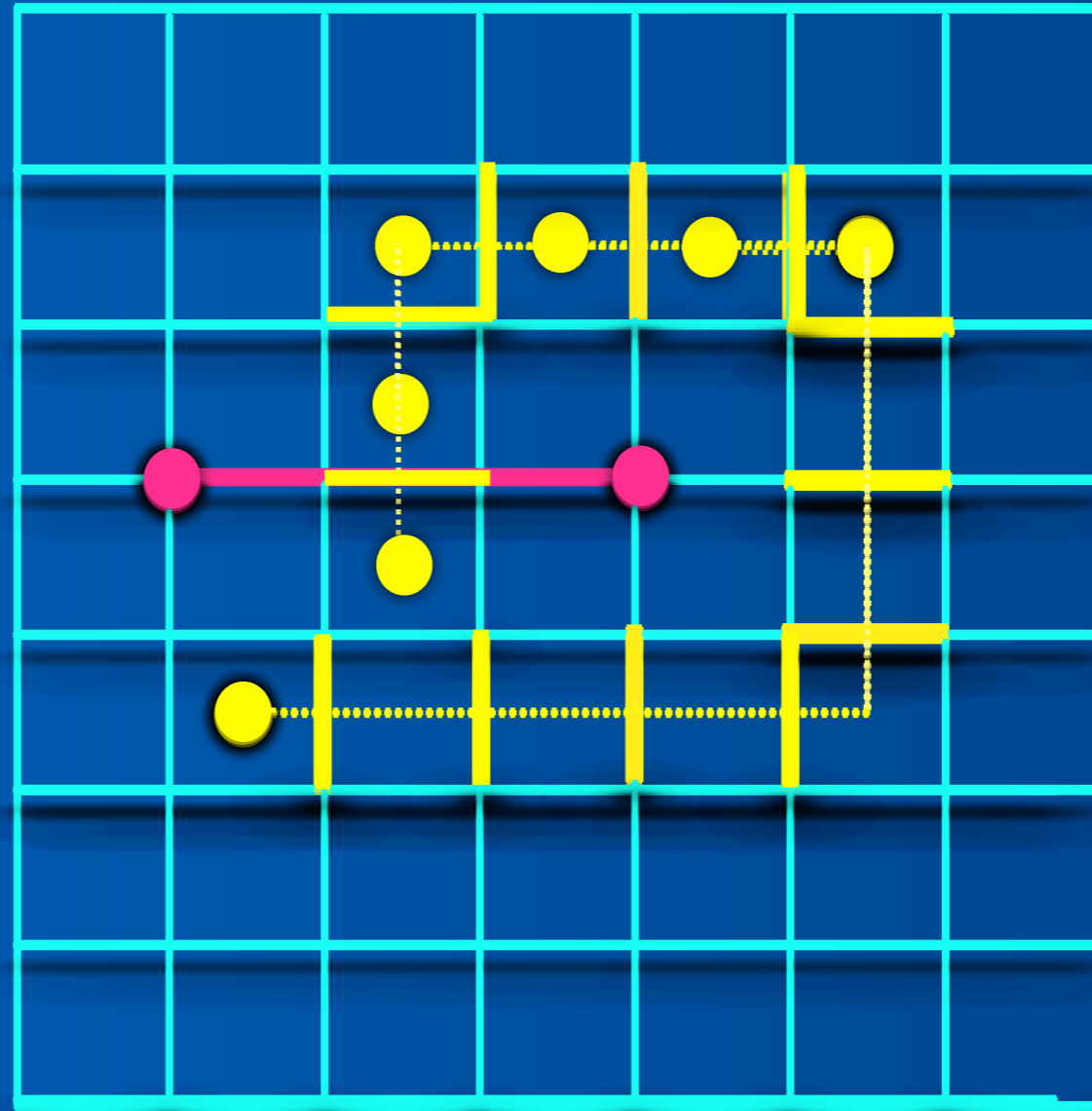
We can implement a Z gate on either of the qubits.

Electric excitations behave as Bosons with respect to each other.

Magnetic excitations behave as Bosons with respect to each other.

But Electric and Magnetic excitations behave as Fermions
with respect to each other.

Why These are Anyons?



So we can do simple,

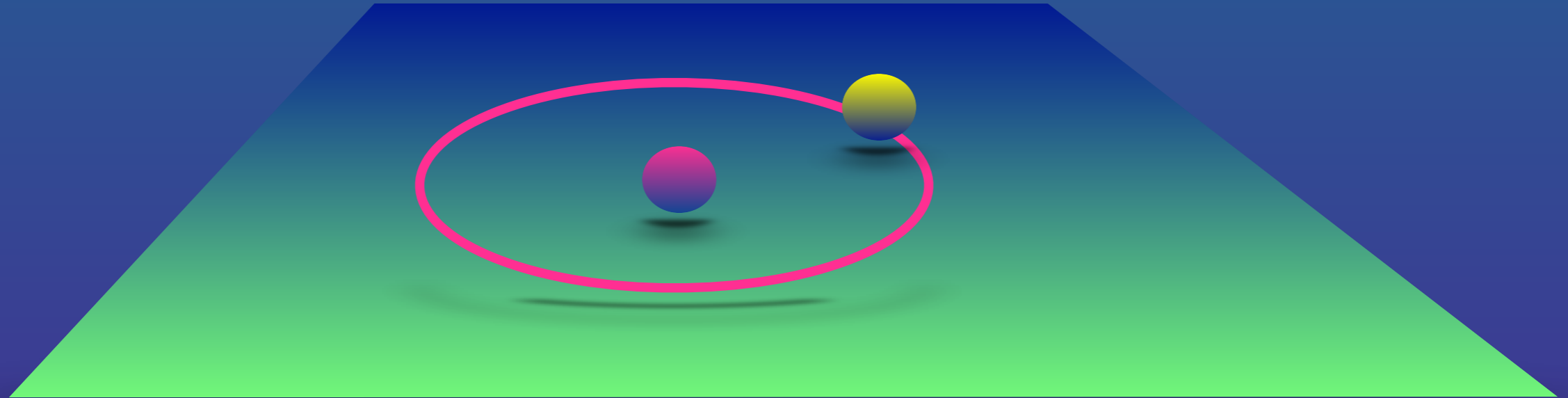
X, Z and Y gates

in a fault-tolerant way.

Unfortunately
the Abelian Models
are not Universal.

We have to consider
Non-Abelian Models.

Non-Abelian Anyons



$$|\psi_i\rangle \rightarrow U_{ij} |\psi_j\rangle$$

Thank you for your attention